# AST/RO ${ }^{13} \mathrm{CO}(J=2 \rightarrow 1)$ AND ${ }^{12} \mathrm{CO}(J=4 \rightarrow 3)$ MAPPING OF SOUTHERN SPITZER c2d SMALL CLOUDS AND CORES 

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#### Abstract

Forty molecular cloud cores in the southern hemisphere from the initial Spitzer Space Telescope Cores-to-Disks (c2d) Legacy program source list have been surveyed in ${ }^{13} \mathrm{CO}(2 \rightarrow 1),{ }^{12} \mathrm{CO}(4 \rightarrow 3)$, and ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ with the Antarctic Submillimeter Telescope and Remote Observatory (AST/RO). The cores, 10 of which contain embedded sources, are located mostly in the Vela, Ophiuchus, Lupus, Chamaeleon, Musca, and Scorpius complexes. ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ emission was undetected in all 40 clouds. We present data of 40 sources in ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ and ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$, significant upper limits of ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$, as well as a statistical analysis of the observed properties of the clouds. We find the typical ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ line width to be $2.0 \mathrm{~km} \mathrm{~s}^{-1}$ for cores with embedded stars and $1.8 \mathrm{~km} \mathrm{~s}^{-1}$ for all others. The typical ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ line width is $2.6-3.7 \mathrm{~km} \mathrm{~s}^{-1}$ for cores with known embedded sources, and $1.6-2.3 \mathrm{~km} \mathrm{~s}^{-1}$ for all others. The average ${ }^{13} \mathrm{CO}$ column density derived from the line intensities was found to be $1.9 \times 10^{15} \mathrm{~cm}^{-2}$ for cores with embedded stars and $1.5 \times 10^{15} \mathrm{~cm}^{-2}$ for all others. The average kinetic temperature in the molecular cores, determined through a large velocity gradient analysis of a set of nine cores, has an average lower limit of 16 K and an average upper limit of 26 K . The average molecular hydrogen density has an average lower limit of $10^{2.9} \mathrm{~cm}^{-3}$ and an average upper limit of $10^{3.3} \mathrm{~cm}^{-3}$ for all cores. For a different subset of nine cores, we have derived masses. They range from 4 to $255 M_{\odot}$. Overall, our c2d sample of southern molecular cores has a range of properties (line width, column density, size, mass, embedded stars) similar to those of past studies.


Subject headings: ISM: clouds - ISM: molecules - radiative transfer - radio lines: ISM — stars: formation submillimeter

## 1. INTRODUCTION

Star formation occurs in dense molecular clouds and their dense cores are the progenitors of protostars (di Francesco et al. 2006; Ward-Thompson et al. 2006). Such cores are compact areas $(<0.5 \mathrm{pc})$ of molecular gas at relatively high density $\left(>10^{3} \mathrm{~cm}^{-3}\right)$. Large efforts have been made in the observation and characterization of such cores. One of the biggest and most recent of these efforts is the Spitzer Space Telescope program "From Molecular Cores to Planet-forming Disks" or "c2d" (Evans et al. 2003). The c2d program surveyed at infrared wavelengths over 100 relatively isolated cores that span the evolutionary sequence from starless cores to protoplanetary disks (Huard et al. 2006). These sources include a wide range of cloud masses and star-forming environments. In order to understand how cores are formed and subsequently evolve into stars, it is necessary to understand the gas dynamics of the underlying molecular cloud. The primary means of studying molecular gas are observations of emission by the rotational transitions of the ground vibrational state of carbon monoxide (CO). These spectral lines occur at frequencies of $J \times$ 115 GHz , for transitions from the $J$ to the $J-1$ state. A realistic picture of the thermodynamic state of a molecular gas can be gained through observations of several transitions of CO isotopes. If the brightness of several of these lines is known, models of the radiative transfer, such as a large velocity gradient (LVG) model (Goldreich \& Kwan 1974), can be used to determine density and temperature of the gas. In order to avoid degeneracy in this solution, it is advisable to have observations of a high- $J$ transition line that is sufficiently high in energy to be only weakly populated.

[^0]From the initial c2d source list of Evans et al. (2003), we have observed with AST/RO all 40 molecular cores that are visible from the South Pole $\left(\delta<-16^{\circ}\right)$, in the ${ }^{13} \mathrm{CO}(2 \rightarrow 1),{ }^{12} \mathrm{CO}(4 \rightarrow 3)$, and ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ transitions. (The initial Spitzer c2d source list was later shortened due to the time available to the c2d program. Some of the sources in this paper have not been observed with Spitzer.) Ten out of the 40 cores are known to have sources embedded in them according to $\operatorname{IRAS}$ (Evans et al. 2003; Lee \& Myers 1999). The detection threshold for star formation using IRAS is typically taken to be $\sim 0.1 L_{\odot}(d / 140 \mathrm{pc})$ (Myers et al. 1987). Cores in which IRAS has not detected a source, which in this paper are referred to as "starless," may still contain faint sources (Young et al. 2004; Bourke et al. 2005, 2006; Huard et al. 2006; Dunham et al. 2006). The c2d team is producing a list of candidate very low-luminosity objects (VeLLOs) in all cores. The multiple transitions observed allow us to perform an LVG analysis and gain information about the kinematic and thermodynamic conditions in a large sample of molecular cores. These observations and results can be used in conjunction with the c2d infrared results for further studies of these cores.

In § 2, we describe the observatory, the observations, and the method of data reduction. Section 3 presents the data set in the form of spectra, maps, tables, and a statistical analysis. The LVG analysis of the radiative transfer is presented in $\S 3.1$, which allows a determination of kinetic temperature and molecular hydrogen density. Masses for several cores are derived in $\S 3.2$. In $\S 4$ we summarize our results and present our conclusions.

## 2. OBSERVATIONS

The observations were performed during the austral winter of 2005 at the Antarctic Submillimeter Telescope and Remote Observatory (AST/RO), which was located at 2847 m elevation at the Amundsen-Scott South Pole Station. These are among the

TABLE 1
Source List

| Source Name | Other Names | $\alpha(\mathrm{J} 2000.0)^{\text {a }}$ | $\delta(\mathrm{J} 2000.0)^{\mathrm{a}}$ | $\begin{gathered} { }^{13} \mathrm{CO}(J=2 \rightarrow 1) \\ \alpha\left({ }^{\prime}\right) \times \delta\left({ }^{\prime}\right) \end{gathered}$ | $\begin{gathered} { }^{12} \mathrm{CO}(J=4 \rightarrow 3) \\ \alpha\left({ }^{\prime}\right) \times \delta\left({ }^{\prime}\right) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| DC 259.5-16.4.................... | BHR 22 | 071412.2 | -48 2910 | $9 \times 9$ | $9 \times 9$ |
| DC 253.8-10.9.................... | BHR 14 | 072932.7 | -411030 | $9 \times 9$ | $9 \times 9$ |
| DC 255.4-3.9..................... | BHR 16 | 080526.3 | -39 0856 | $9 \times 9$ | $9 \times 9$ |
| CG 30, CG 31A-C*b........... | BHR 12, BHR 9-11 | 080904.0 | -36 0051 | $30 \times 30$ | $24.5 \times 23$ |
| DC 257.3-2.5..................... |  | 081709.3 | -39 5225 | $9 \times 15$ | $9.5 \times 15$ |
| DC 266.0-7.5..................... |  | 082122.9 | -49 5047 | $9 \times 9$ | $9 \times 9$ |
| DC 267.4-7.5* .................... | BHR 36 | 082548.3 | -510118 | $12 \times 15$ | $11.5 \times 12$ |
| DC 267.5-7.4... | BHR 37 | 082637.0 | -510231 | $9 \times 9$ | $9 \times 9$ |
| DC 253.6+2.9 | BHR 13 | 082844.0 | -33 4512 | $15 \times 15$ | $14.5 \times 15$ |
| DC 267.6-6.4. | BHR 40 | 083200.8 | -50 3251 | $15 \times 15$ | $9 \times 9$ |
| DC 269.4+3.0. | BHR 43 | 092222.0 | -45 5025 | $12 \times 9$ | $12.5 \times 9$ |
| DC 274.2-0.4..................... | BHR 53 | 092847.0 | -513640 | $9 \times 9$ | $7 \times 7.5$ |
| DC 272.5+2.0 ...................... | BHR 47 | 093102.8 | -483822 | $9 \times 9$ | $11 \times 6.5$ |
| DC 275.9+1.9*..................... | BHR 55 | 094646.0 | -5105 12 | $9 \times 9$ | $9.5 \times 6.5$ |
| DC 291.0-3.5...................... |  | 105943.0 | -63 4343 | $15 \times 15$ | $15 \times 15$ |
| DC 291.1-1.7..................... | BHR 59 | 110652.0 | -62 0643 | $9 \times 9$ | $9 \times 9$ |
| DC 297.7-2.8* .................... | BHR 71 | 120135.9 | -65 0837 | $15 \times 15$ | c |
| DC 300.2-3.5..................... | BHR 75 | 122419.0 | -66 1152 | $9 \times 9$ | $11.5 \times 9$ |
| Mu8. | Musca Dark Cloud | 122935.0 | -71 0935 | $15 \times 21$ | $15 \times 21$ |
| DC 300.7-1.0..................... | BHR 77 | 123134.0 | -63 4432 | $15 \times 15$ | $15 \times 15$ |
| DC 301.2-0.4. | BHR 78 | 123615.0 | -63 1139 | $15 \times 15$ | $9 \times 9$ |
| DC 302.6-15.9.................... | BHR 84 | 124457.6 | -784821 | $15 \times 12$ | $15 \times 9.5$ |
| DC 302.1+7.4*.................... | BHR 83 | 124538.5 | -55 2518 | $9 \times 9$ | $7.5 \times 9$ |
| DC 303.8-14.2*.................. | BHR 86 | 130735.7 | -7700 05 | $15 \times 15$ | $12 \times 15$ |
| DC 327.2+1.8 ...................... | BHR 111 | 154219.2 | -52 4809 | $9 \times 9$ | $9 \times 9$ |
| DC 339.1+11.7 ................... | BHR 129 | 155903.6 | -37 3657 | $15 \times 12$ | $15 \times 12.5$ |
| DC 338.6+9.5 ..................... | BHR 126 | 160428.0 | -39 3756 | $9 \times 9$ | $9.5 \times 7.5$ |
| DC 337.6+7.6 ..................... |  | 160821.0 | -414219 | $18 \times 18$ | $17.5 \times 18$ |
| L43* ................................... | RNO 91 | 163436.2 | -15 4657 | $15 \times 15$ | $15 \times 15$ |
| DC 346.0+7.8 ...................... | B231 | 163654.7 | -35 3608 | $9 \times 12$ | $9.5 \times 12$ |
| DC 346.3+7.8 ..................... |  | 163744.9 | -35 2753 | $9 \times 9$ | $9.5 \times 7.5$ |
| DC 346.4+7.9 ...................... | BHR 144 | 163733.0 | -35 1411 | $9 \times 9$ | $9 \times 9$ |
| L63 .................................... |  | 165011.0 | -18 0434 | $15 \times 15$ | $15 \times 12.5$ |
| CB 68* ............................... | L146 | 165720.5 | -16 0902 | $15 \times 12$ | $15 \times 12$ |
| B59* | L1746 | 171122.1 | -27 2428 | $21 \times 18$ | $20.5 \times 18$ |
| L100* | B62 | 171601.0 | -20 5721 | $15 \times 15$ | $14.5 \times 12$ |
| L1772 ................................. | B65 | 171932.7 | -26 4440 | $15 \times 15$ | $15 \times 15$ |
| B72.................................... | CB 83 | 172345.0 | -23 4128 | $18 \times 12$ | $17.5 \times 12.5$ |
| L328 .................................. | CB 131 | 181700.0 | -180154 | $9 \times 9$ | $8 \times 7.5$ |

[^1]final observations of AST/RO prior to its decommissioning in 2005 December. Low water vapor and high atmospheric stability make this site exceptionally good for submillimeter astronomy (Chamberlin et al. 1997; Lane 1998). AST/RO is a 1.7 m diameter, offset Gregorian telescope that observes at wavelengths between 1.3 mm and $200 \mu \mathrm{~m}$ (Stark et al. 2001; Oberst et al. 2006). Emission from the $220.399 \mathrm{GHz} J=2 \rightarrow 1$ transition of ${ }^{13} \mathrm{CO}$ was mapped using an SIS receiver with $80-160 \mathrm{~K}$ double-sideband receiver noise temperature (Kooi et al. 1992). A dual-channel SIS waveguide receiver (Walker et al. 1992; Honingh et al. 1997) was used for simultaneous observation of 461 GHz and 807 GHz . Single-sideband receiver noise temperatures were $340-440 \mathrm{~K}$ during observations of the $461.041 \mathrm{GHz} J=4 \rightarrow 3$ transition of ${ }^{12} \mathrm{CO}$ and $700-1100 \mathrm{~K}$ for the $J=7 \rightarrow 6$ transition of ${ }^{12} \mathrm{CO}$ at 806.652 GHz .

A multiple-position-switching mode was used with emissionfree reference positions at least $30^{\prime}$ from the mapped area. Inte-
gration times per point were 347 s at 220 GHz and 220 s at $460 / 807 \mathrm{GHz}$. The beam sizes (FWHM) were $180^{\prime \prime}$ at 220 GHz , $103^{\prime \prime}$ at 460 GHz , and $58^{\prime \prime}$ at 807 GHz . All maps were fully sampled with neighboring points $1.5^{\prime}$ apart at 220 GHz and $0.5^{\prime}$ apart at $461 / 807 \mathrm{GHz}$. Tests of the AST/RO pointing model indicate that pointing errors at the time these data were taken could in some cases be as large as $2^{\prime}$ at 220 GHz and $1^{\prime}$ at $461 / 807 \mathrm{GHz}$.

The telescope efficiency $\eta_{\ell}$ was estimated through the reduction of a series of sky dips to $\approx 81 \%$ for the 220 GHz receiver, $\approx 72 \%$ at 461 GHz , and $\approx 80 \%$ at 807 GHz . As described in detail in Stark et al. (2001), the AST/RO main beam efficiency and telescope efficiency are identical to the level that we can measure. The telescope efficiency has been applied to the data as presented below. Atmosphere-corrected system temperatures at the three frequencies ranged from 280 to $480 \mathrm{~K}, 1000$ to 2000 K , and 4000 to $10,000 \mathrm{~K}$, respectively. Two acousto-optical spectrometers (AOS) (Schieder et al. 1989) with 670 kHz channel width, 1.1 MHz


Fig. 1.- ${ }^{13} \mathrm{CO}(2 \rightarrow 1)($ left $)$ and ${ }^{12} \mathrm{CO}(4 \rightarrow 3)($ right $)$ spectra with the highest integrated intensities. Coordinates of the spectra are given in Tables 2 and 3 .
resolution bandwidth, and 1.5 GHz total bandwidth were used as a back end, resulting in velocity resolutions of $0.91 \mathrm{~km} \mathrm{~s}^{-1}$ at $220 \mathrm{GHz}, 0.65 \mathrm{~km} \mathrm{~s}^{-1}$ at 461 GHz , and $0.37 \mathrm{~km} \mathrm{~s}^{-1}$ at 807 GHz . For observations at ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ we used a higher resolution AOS (HRAOS) (Schieder et al. 1989) in addition to the two mentioned above. The HRAOS has a channel width of 31.6 kHz and a resolution bandwidth of 64 MHz , resulting in a velocity resolution of $0.043 \mathrm{~km} \mathrm{~s}^{-1}$ at 220 GHz . The data from this spectrometer are used to derive line widths when the lower resolution was not sufficient and to confirm line widths close to the resolution limit at 220 GHz . These data are also marked in the tables.

A chopper wheel calibration technique was used to observe the sky and a blackbody load of known temperature (Stark et al. 2001) every few minutes during observations. Sky dips were performed frequently to determine the atmospheric transmission. To monitor the pointing accuracy, bright sources (NGC 3576 and/or the moon) were mapped every few hours. After regular intervals or after changes to the receiver systems (tuning), the receiver was manually calibrated against a liquid nitrogen load and a blackbody load of known temperature. This process also corrects for the dark current of the AOS's optical CCDs. The internal consistency of the AST/RO flux scale was estimated from a large-scale map of Lupus, also taken in 2005 with identical instrumental settings. This large map was constructed from smaller maps, with some overlap between them, yielding over a thousand points toward which more than one spectrum was measured at different times. From these data, the flux scale is reproducible at the $30 \%$ level after correction for atmospheric absorption (N. F. H. Tothill et al. 2007, in preparation).

The data taken for this survey were frequency calibrated in several ways. Before and during observations, we observed a comb
spectrum with every automatic receiver calibration to monitor the frequency scale over the entire observing period. During measurements of the sky performed with every automated calibration, the telescope observes a mesospheric telluric line. The velocity scale of all observations was referenced to this scale, resulting in an accuracy of plus/minus one channel.

The data were reduced using the COMB data reduction package as described in Stark et al. (2001). Linear baselines were removed in all spectra at all frequencies, excluding regions where emission was expected based on previous ${ }^{13} \mathrm{CO}(J=1 \rightarrow 0)$ observations (R. Otrupcek 1997, private communication; see also Otrupcek et al. 2000). Removal of first-order baselines proved to be generally sufficient for the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ and ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ data. In very rare cases $(0.01 \%$ of all data), two or three Fourier components had to be removed in addition to the linear baseline. At ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ we did not detect significant $(>3 \sigma)$ emission. However, we were able to deduce upper limits for the ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ intensity from our data for all of the observed sources. Since we detected strong ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ emission in our pointing sources, we can be certain that the nondetection of this line in the cores is not a result of receiver failure. We conclude that the ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ emission from the observed cores is so faint that it lies below the detection sensitivity of our detectors. We thus deduce significant upper limits for this emission.

A list of observed sources in order of increasing R.A., coordinates of the map centers, as well as map sizes is given in Table 1. Cores with known embedded $\operatorname{IRAS}$ sources are marked with asterisks ( ${ }^{*}$ ). The molecular cores we observed contain all cores from the source list of the Spitzer c2d project (Evans et al. 2003) that are visible from the South Pole $\left(\delta<-16^{\circ}\right)$. The observed


Fig. 1-Continued
map sizes were chosen to cover at least the areas observed by c2d but are larger in almost all cases. ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ mapping was sometimes confined to regions where previously mapped ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ emission was observed.

## 3. RESULTS AND ANALYSIS

For each observed core, we present the highest intensity spectrum of both detected transitions in Figures $1 a-1 e$. Observed and derived parameters and source information are listed in Table 2 for the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ transition and in Table 3 for the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ transition. Molecular cores with known embedded IRAS sources are marked with asterisks $\left(^{*}\right)$. The coordinates in Tables 2 and 3 represent the position of the spectrum with the maximum integrated intensity in each map. The source velocity $v_{\text {LSR }}$ was obtained by fitting a Gaussian to the line with maximum intensity. The quantity $\Delta v$ corresponds to the FWHM of this fit. In cases where the line width is lower or close to the spectrometer's velocity resolution of $0.91 \mathrm{~km} \mathrm{~s}^{-1}$, we confirmed the line width from our higher resolution data $\left(0.043 \mathrm{~km} \mathrm{~s}^{-1}\right)$. Cores for which we did this are marked in the table. $\int T_{\mathrm{MB}}(v) d v$ gives the intensity integrated over the brightest line in the map, and rms is the root mean squared noise of the channels in the spectrum.

In Table 2, the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ intensity is converted to a total column density of ${ }^{13} \mathrm{CO}\left(N_{13} \mathrm{CO}\right)$ by assuming local thermodynamic equilibrium (LTE):

$$
\begin{equation*}
N^{{ }^{13}} \mathrm{CO}=1.51 \times 10^{14} \frac{T_{\mathrm{ex}} e^{5.3 / T} \int \tau(v) d v}{1-e^{-10.6 / T_{\mathrm{ex}}}} \tag{1}
\end{equation*}
$$

Assuming optically thin ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ emission, that the source fills the main beam, that the background temperature of 2.7 K can be neglected and that the Rayleigh-Jeans correction can be neglected as well, we use the relation

$$
\begin{equation*}
T_{\mathrm{ex}} \tau(v)=T_{\mathrm{MB}}(v) \tag{2}
\end{equation*}
$$

where $\tau(v)$ is the optical depth of the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ line, $T_{\mathrm{MB}}$ is the main beam brightness temperature, and $T_{\text {ex }}$ is the excitation temperature, which is usually taken to be the kinetic temperature in the cloud: $T_{\mathrm{ex}}=T_{\mathrm{kin}}=10 \mathrm{~K}$. (Up to $T_{\mathrm{ex}}=40 \mathrm{~K}$ the effect of this parameter on the ${ }^{13} \mathrm{CO}$ column density is less than a factor of 2.)

In Figure 2, we show histograms of the line widths of the two detected transitions. The distribution of the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ line width has a mean of $1.8 \mathrm{~km} \mathrm{~s}^{-1}$ and a standard deviation of $0.4 \mathrm{~km} \mathrm{~s}^{-1}$. The cores with embedded sources are represented by the black bars. It is noticeable that their distribution is generally skewed to higher values. The mean of their line widths distribution is $2.0 \mathrm{~km} \mathrm{~s}^{-1}$ and the standard deviation is $0.4 \mathrm{~km} \mathrm{~s}^{-1}$. Cores with embedded stars tend to be warmer and more turbulent, and their emission lines are therefore subject to more broadening than are cores without embedded stars. If the gas kinetic temperature is typically 20 K , the corresponding optically thin thermal line width of CO is only $0.2 \mathrm{~km} \mathrm{~s}^{-1}$, so the observed widths are due primarily to turbulent motions and optical depth effects. DC 253.8-10.9 shows an exceptionally narrow line with only $0.8 \mathrm{~km} \mathrm{~s}^{-1}$. It also is the weakest core in ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$. L328 has the broadest line of the cores with no embedded $I R A S$ source at $2.7 \mathrm{~km} \mathrm{~s}^{-1}$.


Fig. 1-Continued
This core has recently been found to have a VeLLO embedded (C. W. Lee 2007, in preparation). The widest line at $2.9 \mathrm{~km} \mathrm{~s}^{-1}$ is emitted by DC $275.9+1.9$, which does have an embedded star (Bourke et al. 1995).

The distribution of the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ line width also shows that the lines from cores with embedded sources tend to be broader than lines from starless cores. The mean line width of cores with embedded stars is $3.2 \mathrm{~km} \mathrm{~s}^{-1}$, with a standard deviation of $1.3 \mathrm{~km} \mathrm{~s}^{-1}$. Cores with no embedded stars have a mean line width of $2.0 \mathrm{~km} \mathrm{~s}^{-1}$, with a standard deviation of $0.7 \mathrm{~km} \mathrm{~s}^{-1}$. For this histogram, we did not consider cores whose peak intensity is lower than $5 \sigma$. These omitted cores are DC 302.6-15.9, L63, CB 68, and L328. The widest line of all cores is found in DC 303.8-14.2 at $5.8 \mathrm{~km} \mathrm{~s}^{-1}$, which has an embedded star (Bourke et al. 1995). The widest line of cores with no embedded sources is emitted by DC 300.7-1.0 with $3.5 \mathrm{~km} \mathrm{~s}^{-1}$.

The survey of southern cores by Vilas-Boas et al. (1994) in ${ }^{13} \mathrm{CO}(1 \rightarrow 0)$ reports an average line width of $0.8 \mathrm{~km} \mathrm{~s}^{-1}$. The only source observed in both the Vilas-Boas et al. survey and this survey is CG 30 . Vilas-Boas et al. (1994) gives a line width of $1.19 \mathrm{~km} \mathrm{~s}^{-1}$ obtained with a resolution of $0.12 \mathrm{~km} \mathrm{~s}^{-1}$. Using our high-resolution spectrometer, we observed a line width of $1.2 \pm$ $0.1 \mathrm{~km} \mathrm{~s}^{-1}$, which is in good agreement with Vilas-Boas et al. (1994).

A histogram of the ${ }^{13} \mathrm{CO}$ column density is shown in Figure 3. The mean ${ }^{13} \mathrm{CO}$ column density for cores with no embedded stars is $(0.5-2.0) \times 10^{15} \mathrm{~cm}^{-2}$ and slightly higher, around $(1.0-2.0) \times$ $10^{15} \mathrm{~cm}^{-2}$, for cores with embedded stars, which do not show ${ }^{13} \mathrm{CO}$ column densities lower than $1.0 \times 10^{15} \mathrm{~cm}^{-2}$. Fifty-five percent of the cores have a ${ }^{13} \mathrm{CO}$ column density between $0.5 \times 10^{15}$
and $1.5 \times 10^{15} \mathrm{~cm}^{-2}$, and $75 \%$ have column densities between $0.5 \times 10^{15}$ and $2.0 \times 10^{15} \mathrm{~cm}^{-2}$. Assuming a ${ }^{13} \mathrm{CO} / \mathrm{H}_{2}$ ratio of $1.7 \times 10^{-6}$ (Wilson \& Rood 1994; Frerking et al. 1982), we can convert the ${ }^{13} \mathrm{CO}$ column densities to the $\mathrm{H}_{2}$ column density. The $\mathrm{H}_{2}$ column density distribution has a mean of $1.9 \times 10^{21} \mathrm{~cm}^{-2}$ and a standard deviation of $0.7 \times 10^{21} \mathrm{~cm}^{-2}$ for cores with embedded sources. For all other sources, the mean is $1.5 \times 10^{21} \mathrm{~cm}^{-2}$ and the standard deviation is $0.7 \times 10^{21} \mathrm{~cm}^{-2}$. Cores with embedded sources tend to have higher column densities. Seven of the observed cores lie in Ophiuchus, which has a high level of star formation activity (Wilking 1992). In these cores the $\mathrm{H}_{2}$ column density assumes high values, with a mean of $3.7 \times 10^{21} \mathrm{~cm}^{-2}$ and a dispersion of $1.5 \times 10^{21} \mathrm{~cm}^{-2}$. Twenty-three cores lie within clouds where the star-forming activity is low (Vela, Musca, Coalsack, Chamaeleon II and III, Lupus, Scorpius). These cores have a mean $\mathrm{H}_{2}$ column density of $2.7 \times 10^{21} \mathrm{~cm}^{-2}$ with a dispersion of $1.0 \times 10^{21} \mathrm{~cm}^{-2}$. This difference as an indicator of star-forming activity was also observed by Vilas-Boas et al. (2000) in a ${ }^{13} \mathrm{CO}(1-0)$ survey of different (none of the sources in their and our survey overlap) dense condensations in dark clouds. Vilas-Boas et al. (2000) set the limit for star formation-enabling $\mathrm{H}_{2}$ column density to as high as $1.1 \times 10^{22} \mathrm{~cm}^{-2}$. This limit appears to be too high, considering our results for B59. B59 in Ophiuchus is the strongest core observed and has the exceptionally high ${ }^{13} \mathrm{CO}$ column density of $4.3 \times 10^{15} \mathrm{~cm}^{-2}$ and $\mathrm{H}_{2}$ column density of $7.3 \times 10^{21} \mathrm{~cm}^{-2}$, thus indicating strong star formation activity. A group of embedded sources has already been found in this core (Brooke et al. 2007).

### 3.1. Large Velocity Gradient Analysis

For a more detailed analysis of our data, we chose only the cores most suitable for the purpose: cores with strong ( $>5 \sigma$ ) emission in both ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ and ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$. We also made sure there was no large pointing offset between intensity peaks at the two wavelengths. In Figure 4, we present maps of the nine cores chosen for the LVG analysis. The AST/RO contour maps are superimposed on optical images from the Digital Sky Survey (DSS) retrieved from SkyView (McGlynn et al. 1996). The cores are shown in order of ascending R.A. We present the majority of cores superimposed on $15^{\prime} \times 15^{\prime}$ images, but AST/RO map sizes vary. Maps in the same row show the same source, the left column shows the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ emission, the right column shows the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$. All maps are labeled with the source name, the observed transition, and the maximum integrated intensity in $\mathrm{K} \mathrm{km} \mathrm{s}^{-1}$. Contours in all maps are in percentage of the peak intensities. We show contour levels in steps of $10 \%$ of the highest contour, whose value is given in each figure. The first maps on each page show the AST/RO effective beam sizes of $216^{\prime \prime}$ at 220 GHz and $117^{\prime \prime}$ at 461 GHz . Since the maps were made using Gaussian interpolation, the effective beam sizes in the maps are slightly larger than the actual telescope beam sizes.

Maps that are larger than $15^{\prime} \times 15^{\prime}$ are shown in order of ascending R.A. after the set of $15^{\prime} \times 15^{\prime}$ maps. The left panels show the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ map, the right panels show the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ map. In cases with more than one source per map, the sources are labeled individually.

We can estimate the kinetic temperature $T_{\text {kin }}$, and the number density of molecular hydrogen $n\left(\mathrm{H}_{2}\right)$, using a LVG analysis of the radiative transfer (Goldreich \& Kwan 1974). Our LVG radiative transfer code has been developed and applied by M. Yan and S. Kim (Kim et al. 2002) and has been implemented as a java applet. ${ }^{3}$ This model simulates a plane-parallel cloud geometry and

[^2]TABLE 2
Observed and Derived Parameters of the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ Lines

| Source | $\alpha(\mathrm{J} 2000.0)^{\text {a }}$ | $\delta(\mathrm{J} 2000.0)^{\mathrm{a}}$ | $\begin{gathered} v_{\mathrm{LSR}} \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} \Delta v \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{aligned} & \int T_{\mathrm{MB}}(v) d v \\ & \left(\mathrm{~K} \mathrm{~km} \mathrm{~s}^{-1}\right) \end{aligned}$ | rms Noise (K) | $\begin{gathered} N_{13} \mathrm{CO} \\ \left(10^{15} \mathrm{~cm}^{-2}\right) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| DC 259.5-16.4................... | 071403.2 | -48 2740 | 4.6 | $1.5 \pm 0.1$ | $2.7 \pm 0.2$ | 0.1 | 1.1 |
| DC 253.8-10.9. | 072932.7 | -410730 | -0.6 | $0.8 \pm 0.1^{\text {b }}$ | $1.5 \pm 0.2$ | 0.2 | 0.6 |
| DC 255.4-3.9..................... | 080526.3 | -39 0726 | 9.8 | $2.4 \pm 0.1$ | $2.2 \pm 0.2$ | 0.2 | 0.9 |
| CG 31A-C | 080841.8 | -35 5451 | 6.7 | $1.9 \pm 0.1$ | $6.0 \pm 0.2$ | 0.2 | 2.4 |
| CG 30**. | 080926.3 | -360221 | 6.5 | $1.2 \pm 0.1^{\text {b }}$ | $4.0 \pm 0.2$ | 0.1 | 1.6 |
| DC 257.3-2.5. | 081717.1 | -39 4755 | 8.8 | $2.0 \pm 0.2$ | $6.8 \pm 0.5$ | 0.6 | 2.7 |
| DC 266.0-7.5.. | 082122.9 | -49 4747 | 4.4 | $1.7 \pm 0.1$ | $4.8 \pm 0.2$ | 0.1 | 1.9 |
| DC 267.4-7.5*. | 082548.3 | -50 5948 | 5.8 | $1.7 \pm 0.1$ | $5.3 \pm 0.2$ | 0.2 | 2.1 |
| DC 267.5-7.4.. | 082646.5 | -5104 01 | 6.5 | $1.5 \pm 0.1$ | $2.8 \pm 0.2$ | 0.2 | 2.9 |
| DC 253.6+2.9 ...................... | 082844.0 | -33 4342 | 7.3 | $1.7 \pm 0.1$ | $4.7 \pm 0.2$ | 0.2 | 1.9 |
| DC 267.6-6.4. | 083210.2 | -50 3121 | 6.3 | $1.6 \pm 0.1$ | $2.8 \pm 0.2$ | 0.1 | 1.1 |
| DC 269.4+3.0 | 092222.0 | -45 5455 | -2.9 | $1.3 \pm 0.2^{\text {b }}$ | $1.9 \pm 0.1$ | 0.1 | 0.9 |
| DC 274.2-0.4. | 092847.0 | -513510 | 6.4 | $2.2 \pm 0.2$ | $2.5 \pm 0.2$ | 0.2 | 1.0 |
| DC 272.5+2.0 | 093102.8 | -48 3652 | -3.9 | $1.6 \pm 0.2$ | $1.6 \pm 0.2$ | 0.1 | 0.6 |
| DC 275.9+1.9*. | 094636.5 | $-510042$ | -4.9 | $2.9 \pm 0.2$ | $3.4 \pm 0.2$ | 0.1 | 1.3 |
| DC 291.0-3.5.. | 110010.1 | -63 4043 | -3.8 | $1.7 \pm 0.1$ | $2.9 \pm 0.2$ | 0.1 | 1.1 |
| DC 291.1-1.7. | 110704.8 | -62 0343 | -4.2 | $1.8 \pm 0.1$ | $3.2 \pm 0.2$ | 0.1 | 1.3 |
| DC 297.7-2.8*.................... | 120157.0 | -650519 | -3.9 | $2.1 \pm 0.1$ | $7.4 \pm 0.2$ | 0.2 | 2.9 |
| DC 300.2-3.5. | 122503.5 | -66 0852 | -5.2 | $1.9 \pm 0.2$ | $2.3 \pm 0.2$ | 0.2 | 0.9 |
| Mu8. | 123030.5 | -71 0335 | 3.8 | $1.7 \pm 0.1$ | $4.4 \pm 0.2$ | 0.1 | 1.7 |
| DC 300.7-1.0.. | 123120.5 | -63 4132 | -5.5 | $1.1 \pm 0.2^{\text {b }}$ | $2.0 \pm 0.1$ | 0.2 | 0.8 |
| DC 301.2-0.4.. | 123601.7 | -63 0839 | -4.4 | $1.6 \pm 0.1$ | $1.8 \pm 0.2$ | 0.1 | 0.7 |
| DC 302.6-15.9. | 124426.8 | -78 4521 | 4.1 | $1.8 \pm 0.1$ | $1.8 \pm 0.2$ | 0.1 | 0.7 |
| DC 302.1+7.4*. | 124549.1 | -55 2218 | -14.7 | $2.1 \pm 0.2$ | $2.6 \pm 0.2$ | 0.1 | 1.0 |
| DC 303.8-14.2* .................. | 130735.7 | -76 5705 | 4.3 | $2.1 \pm 0.1$ | $3.6 \pm 0.2$ | 0.1 | 1.4 |
| DC 327.2+1.8 ...................... | 154249.0 | -52 4939 | -0.4 | $2.0 \pm 0.2$ | $2.3 \pm 0.2$ | 0.1 | 0.9 |
| DC 339.1+11.7 .................... | 155911.2 | -37 3657 | 5.7 | $1.7 \pm 0.1$ | $4.3 \pm 0.2$ | 0.1 | 1.7 |
| DC 338.6+9.5 | 160443.6 | -39 3626 | 5.2 | $1.0 \pm 0.1^{\text {b }}$ | $3.1 \pm 0.2$ | 0.1 | 1.2 |
| DC 337.6+7.6 | 160821.0 | -414049 | 5.4 | $1.7 \pm 0.1$ | $7.1 \pm 0.1$ | 0.2 | 2.8 |
| L43* | 163423.7 | $-154957$ | 1.0 | $1.8 \pm 0.1$ | $6.6 \pm 0.2$ | 0.1 | 2.6 |
| DC 346.0+7.8 | 163654.7 | -35 3738 | 4.1 | $2.2 \pm 0.1$ | $5.4 \pm 0.2$ | 0.2 | 2.1 |
| DC 346.3+7.8 | 163737.5 | -35 3053 | 3.8 | $1.8 \pm 0.1$ | $4.7 \pm 0.2$ | 0.2 | 1.9 |
| DC 346.4+7.9 . | 163747.7 | -351411 | 4.1 | $1.7 \pm 0.1$ | $4.0 \pm 0.2$ | 0.2 | 1.6 |
| L63. | 164952.1 | -1803 04 | 6.3 | $2.0 \pm 0.1$ | $7.3 \pm 0.1$ | 0.2 | 2.9 |
| CB 68* | 165649.3 | -160432 | 5.1 | $1.7 \pm 0.1$ | $3.2 \pm 0.1$ | 0.2 | 1.3 |
| B59* | 171126.6 | -27 2558 | 3.7 | $2.1 \pm 0.1$ | $10.9 \pm 0.2$ | 0.2 | 4.3 |
| L100* | 171601.0 | -20 5721 | 2.1 | $2.2 \pm 0.1$ | $4.1 \pm 0.3$ | 0.1 | 1.6 |
| L1772 | 171932.7 | -26 4140 | 3.9 | $1.8 \pm 0.1$ | $3.5 \pm 0.2$ | 0.2 | 1.4 |
| B72. | 172338.5 | -23 3828 | 5.4 | $1.7 \pm 0.1$ | $3.4 \pm 0.2$ | 0.2 | 1.3 |
| L328 ..................................... | 181653.7 | -18 0324 | 7.4 | $2.7 \pm 0.2$ | $2.8 \pm 0.2$ | 0.2 | 1.1 |

[^3]uses the CO collisional rates from Turner (1995) newly derived values for the $\mathrm{H}_{2}$ ortho-to-para ratio ( $\approx 2$ ), and the collisional quenching rate of CO by $\mathrm{H}_{2}$ impact (Balakrishnan et al. 2002). The model has two input parameters: the ratio of ${ }^{12} \mathrm{CO}$ to ${ }^{13} \mathrm{CO}$ abundance and the ratio $X(\mathrm{CO}) / \nabla V$, where $X(\mathrm{CO})$ is the fractional CO abundance and $\nabla V$ denotes the velocity gradient. The abundance ratio ${ }^{12} \mathrm{CO} /{ }^{13} \mathrm{CO}$ is taken to be 60 in dense molecular clouds (Wilson \& Rood 1994). The ${ }^{12} \mathrm{CO} / \mathrm{H}_{2}$ ratio is taken to be $10^{-4}$ (Frerking et al. 1982) (leading to a ${ }^{13} \mathrm{CO}$ to $\mathrm{H}_{2}$ ratio of $1.7 \times 10^{-6}$ ) and the velocity gradient across the cores is $1 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{pc}^{-1}$.

The ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ data were convolved with a Gaussian kernel to the lower ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ resolution. The resulting peak intensities are listed in Table 4.

Although we were unable to detect any significant ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ emission, we can use these data to derive a significant upper limit for this emission for each source. We convolved the higher reso-
lution ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ data with a Gaussian kernel to the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ resolution and derived a $3 \sigma$ upper limit for the emission through this relation:

$$
\begin{equation*}
3 \sigma=3 \frac{\mathrm{rms} \times \sqrt{N} \times \text { channel width } \times 0.94}{\operatorname{FWHM}(4-3)} \tag{3}
\end{equation*}
$$

where "rms" is the rms noise of the co-added spectrum; $N$ is the number of channels over which the rms was measured (twice the expected line width of the ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ line $)$; "channel width" is the width of the acousto-optic spectrometer channels at this wavelength; and "FWHM(4-3)" is the half-power width of the corresponding ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ line. The factor of 0.94 comes from the relation for Gaussian profiles: $\mathrm{FWHM} \times$ peak $=0.94 \times$ area.

The results are shown in Table 4, where upper and lower limits are denoted "u.l." and "1.l." Since the line brightness temperature

TABLE 3
Observed Parameters of the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ Lines

| Source | $\alpha(\mathrm{J} 2000.0)^{\mathrm{a}}$ | $\delta(\mathrm{J} 2000.0)^{\mathrm{a}}$ | $\begin{gathered} v_{\mathrm{LSR}} \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} \Delta v \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{aligned} & \int T_{\mathrm{MB}}(v) d v \\ & \left(\mathrm{~K} \mathrm{~km} \mathrm{~s}^{-1}\right) \end{aligned}$ | rms Noise <br> (K) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| DC 259.5-16.4.................. | 071409.2 | -4828 10 | 4.2 | $3.0 \pm 0.4$ | $3.9 \pm 0.4$ | 0.4 |
| DC 253.8-10.9.. | 072932.7 | -410918 | -1.1 | $3.1 \pm 0.3$ | $4.1 \pm 0.3$ | 0.3 |
| DC 255.4-3.9. | 080522.2 | -39 1126 | 9.7 | $2.1 \pm 0.3$ | $3.2 \pm 0.3$ | 0.4 |
| CG 31A-C........................ | 080901.5 | -35 5921 | 6.7 | $2.2 \pm 0.1$ | $14.4 \pm 0.5$ | 0.3 |
| CG $30 *$ * | 080933.7 | -36 0451 | 6.5 | $4.2 \pm 0.3$ | $11.0 \pm 0.3$ | 0.6 |
| DC 257.3-2.5.................... | 081715.7 | -39 5125 | 8.6 | $1.4 \pm 0.1$ | $5.9 \pm 0.4$ | 0.4 |
| DC 266.0-7.5. | 082113.6 | -49 5447 | 3.8 | $2.0 \pm 0.2$ | $4.9 \pm 0.3$ | 0.3 |
| DC 267.4-7.5*. | 082545.1 | -50 5918 | 5.5 | $1.8 \pm 0.1$ | $10.9 \pm 0.2$ | 0.2 |
| DC 267.5-7.4.. | 082648.5 | $-510231$ | 6.0 | $2.7 \pm 0.2$ | $5.9 \pm 0.3$ | 0.3 |
| DC 253.6+2.9 . | 082852.6 | -33 4642 | 6.3 | $1.3 \pm 0.1$ | $6.4 \pm 0.3$ | 0.3 |
| DC 267.6-6.4. | 083210.3 | -5033 39 | 5.8 | $1.6 \pm 0.1$ | $5.2 \pm 0.3$ | 0.3 |
| DC 269.4+3.0 . | 092227.7 | -45 5025 | -1.4 | $2.0 \pm 0.2$ | $3.3 \pm 0.2$ | 0.3 |
| DC 274.2-0.4. | 092850.2 | -513640 | 5.9 | $2.2 \pm 0.2$ | $2.6 \pm 0.2$ | 0.2 |
| DC 272.5+2.0 | 093111.9 | -48 3734 | -3.8 | $2.7 \pm 0.9$ | $1.5 \pm 0.2$ | 0.3 |
| DC 275.9+1.9*. | 094634.0 | $-510112$ | -5.7 | $2.6 \pm 0.3$ | $4.2 \pm 0.4$ | 0.4 |
| DC 291.0-3.5.. | 110010.1 | -63 4043 | -4.8 | $1.0 \pm 0.1$ | $5.0 \pm 0.4$ | 0.6 |
| DC 291.1-1.7. | 110630.6 | -62 0313 | -4.3 | $1.2 \pm 0.1$ | $3.5 \pm 0.2$ | 0.3 |
| DC 300.2-3.5. | 122419.0 | -661252 | -5.7 | $1.1 \pm 0.1$ | $2.7 \pm 0.2$ | 0.2 |
| Mu8. | 122807.6 | -71 1905 | 3.3 | $2.5 \pm 0.3$ | $6.4 \pm 0.5$ | 0.7 |
| DC 300.7-1.0. | 123129.5 | -63 4432 | -4.9 | $3.5 \pm 0.1$ | $5.9 \pm 0.1$ | 1.3 |
| DC 301.2-0.4.................... | 123619.4 | -63 1109 | -4.7 | $2.6 \pm 0.5$ | $2.4 \pm 0.3$ | 0.3 |
| DC 302.6-15.9................... | 124538.7 | -78 4721 | 3.8 | $1.7 \pm 0.1$ | $2.2 \pm 0.1$ | 0.5 |
| DC 302.1+7.4*.................... | 124538.5 | -55 2448 | -14.7 | $2.2 \pm 0.3$ | $3.7 \pm 0.4$ | 0.4 |
| DC 303.8-14.2* ................ | 130726.8 | -77 0135 | 3.4 | $5.8 \pm 0.5$ | $6.7 \pm 0.3$ | 0.2 |
| DC 327.2+1.8 | 154232.4 | -52 4709 | -0.8 | $3.2 \pm 0.5$ | $3.1 \pm 0.3$ | 0.3 |
| DC 339.1+11.7 | 155913.6 | -37 3557 | 5.1 | $1.3 \pm 0.1$ | $6.6 \pm 0.5$ | 0.5 |
| DC 338.6+9.5. | 160433.2 | -39 3856 | 4.4 | $1.0 \pm 0.1$ | $3.6 \pm 0.3$ | 0.3 |
| DC 337.6+7.6 | 160746.2 | -414349 | 4.7 | $1.7 \pm 0.1$ | $9.5 \pm 0.1$ | 0.5 |
| L43* | 163425.8 | -154757 | 0.4 | $3.0 \pm 0.3$ | $14.8 \pm 1.0$ | 1.0 |
| DC 346.0+7.8 | 163657.2 | -35 3638 | 3.3 | $1.8 \pm 0.2$ | $6.8 \pm 0.4$ | 0.4 |
| DC 346.3+7.8 ...................... | 163752.3 | -35 2953 | 3.4 | $1.8 \pm 0.1$ | $7.1 \pm 0.3$ | 0.3 |
| DC 346.4+7.9 ...................... | 163757.5 | -35 1511 | 3.3 | $1.6 \pm 0.1$ | $4.8 \pm 0.2$ | 0.3 |
| L63 ..................................... | 165004.7 | -180104 | 5.8 | $1.2 \pm 0.1$ | $6.7 \pm 0.2$ | 1.4 |
| CB 68* | 165707.9 | $-160532$ | 5.6 | $3.4 \pm 0.4$ | $7.5 \pm 0.4$ | 1.2 |
| B59* | 171117.6 | -27 2528 | 4.0 | $3.6 \pm 0.3$ | $18.1 \pm 0.7$ | 0.9 |
| L100* ............................... | 171609.5 | -20 6621 | 1.4 | $1.8 \pm 0.3$ | $14.2 \pm 1.6$ | 0.4 |
| L1772. | 171950.6 | -26 4340 | 3.5 | $2.1 \pm 0.2$ | $5.1 \pm 0.3$ | 0.4 |
| B72 .................................... | 172344.9 | -23 4028 | 5.1 | $1.9 \pm 0.3$ | $6.0 \pm 0.5$ | 0.8 |
| L328 ................................... | 181702.1 | -1802 42 | 6.9 | $1.2 \pm 0.3$ | $3.3 \pm 0.7$ | 0.9 |

${ }^{\text {a }}$ Coordinates of spectrum with highest integrated intensity. Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.
${ }^{\mathrm{b}}$ Cores with known embedded sources as detected using IRAS are marked with asterisks ( ${ }^{*}$ ).


Fig. 2.-Histograms of the ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ line width (left) and ${ }^{12} \mathrm{CO}(4 \rightarrow 3)($ right $)$. Gray bars represent sources with no embedded stars, black bars represent sources with embedded IRAS sources.


Fig. 3.-Histogram of the ${ }^{13} \mathrm{CO}$ column density for all sources. Gray bars represent sources with no embedded stars, black bars represent sources with embedded IRAS sources.
$T_{7 \rightarrow 6}^{12}$ of the ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ line is an upper limit, the ratio $T_{7 \rightarrow 6}^{12} / T_{4 \rightarrow 3}^{12}$, the kinetic temperature and the $\mathrm{H}_{2}$ density are also upper limits. To obtain the lower temperature and density limits, we assume no ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ emission at all: $T_{7 \rightarrow 6}^{12}=0$.

As can be seen from Table 4, the lower temperature limits average to 16 K and are not greater than 20 K for any source. The distribution of the upper temperature limits have an average of 26 K . The $\mathrm{H}_{2}$ density distribution is fairly uniform, ranging from $10^{2.7} \mathrm{~cm}^{-3}$ to $10^{3.7} \mathrm{~cm}^{-3}$ with an average lower limit of $10^{2.9} \mathrm{~cm}^{-3}$ and an average upper limit of $10^{3.3} \mathrm{~cm}^{-3}$ for all sources. Due to the small sample size, we cannot detect a significant difference in the mean densities of cores with and without embedded stars.

### 3.2. Masses

Nine cores were chosen for the derivation of their masses. We selected cores whose half-peak contours in ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ are closed to derive the cores' masses. These cores are also resolved. Maps of the selected sources are presented in Figures 4 and 5. We show contour levels in steps of $10 \%$ of the highest contour, whose value is given in each figure. DC 302.6-15.9 and DC 303.8-14.2 are at very high elevation ( $79^{\circ}$ and $77^{\circ}$ ) and possibly suffer from a systematic pointing offset as suggested from the background image. For the calculation of their masses, a potential pointing offset is irrelevant. There is no evidence that other cores show a similar pointing offset. We can obtain the total mass of the core using

$$
\begin{equation*}
M=\langle m\rangle \int_{A} N\left(\mathrm{H}_{2}\right) d A \tag{4}
\end{equation*}
$$

where $A$ is the projected area and $\langle m\rangle$ is the mean mass per hydrogen molecule. This mean mass per hydrogen molecule $\langle m\rangle$
accounts for a helium abundance of one helium atom per five hydrogen molecules. We take into account that the sources are not spherical, but elliptical, and list their two radii (major and minor axes) in Table 5 . We derive an average ${ }^{13} \mathrm{CO}$ column density for the core area that is defined by the $50 \%$ peak intensity contour by co-adding all data that were taken in this area and making use of equation (1). This conversion assumes an excitation temperature of $T_{\mathrm{ex}}=10 \mathrm{~K}$. Up to $T_{\mathrm{ex}}=40 \mathrm{~K}$, the effect of this parameter on the ${ }^{13} \mathrm{CO}$ column density is less than a factor of 2 . By assuming a ratio of $1.7 \times 10^{-6}$ between ${ }^{13} \mathrm{CO}$ and $\mathrm{H}_{2}$ (Wilson \& Rood 1994; Frerking et al. 1982), we then convert the ${ }^{13} \mathrm{CO}$ column densities to the $\mathrm{H}_{2}$ column densities listed in Table 5 and use the relation above to determine their masses. Their distances are quoted from the literature: DC 259.5-16.4, CG 30, DC $267.4-7.5$, DC 253.6+2.0, and DC 274.2-0.4 lie within the Vela complex at $450 \pm 50 \mathrm{pc}$ (Woermann et al. 2001); DC 297.7-2.8 is associated with the Coalsack at $150 \pm 30$ pc (Franco 1995; Corradi et al. 1997); DC 302.6-15.9 and DC 303.8-14.2 lie in the Chamaeleon cloud complexes at distances between $150-180 \mathrm{pc}$ (Whittet et al. 1997; Knude \& Høg 1998); and L100 lies within Ophiuchus at $125 \pm 25 \mathrm{pc}$ (de Geus et al. 1989). The results are listed in Table 5. If the upper limit kinetic temperature was obtained through the LVG analysis, we list it again.

The derived masses range from 4 to $255 M_{\odot}$, depending on the size of the core. The highest mass core, DC 267.4-7.5, is also the largest. The smallest and weakest core, DC 302.6-15.9, does have the lowest mass. A correlation between temperature and mass cannot be established.

We have found three cores whose masses have been derived from different observations. Vilas-Boas et al. (1994) derived the mass of CG 30 to be $43 M_{\odot}$ with an uncertainty of a factor of 2.5 , but they also quote a shorter distance of 300 pc . With a distance of 450 pc , this mass would be $97 M_{\odot}$, which is within a factor of 2 of our value ( $58 M_{\odot}$ ). Bourke et al. (1997) have calculated the mass of DC 297.7-2.8 to be $40 M_{\odot}$, using a larger radius and larger distance for the source. With the distance we use, 150 pc , this value becomes $23 M_{\odot}$, which is also within a factor of 2 of our value ( $12 M_{\odot}$ ). Keto \& Myers (1986) calculated the mass of DC $259.5-16.4$ to be $41 M_{\odot}$, but they used the much smaller distance of 100 pc . Combining their observations with a distance of 450 pc leads to a mass for DC $259.5-16.4$ of $830 M_{\odot}$.

## 4. SUMMARY AND CONCLUSIONS

In order to study the physical conditions in small isolated molecular clouds, we have mapped 40 cores in the southern hemisphere in ${ }^{13} \mathrm{CO}(2 \rightarrow 1),{ }^{12} \mathrm{CO}(4 \rightarrow 3)$, and ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ with AST/RO. The sources cover all cores visible from the South Pole $\left(\delta<-16^{\circ}\right)$ from the initial c2d source list (Evans et al. 2003). The AST/RO survey at millimeter and submillimeter wavelengths provides a direct diagnostic of the dominant gas components in the clouds. By observing the ${ }^{12} \mathrm{CO}$ and ${ }^{13} \mathrm{CO}$ lines, we have obtained information about the physical conditions within the cloud. ${ }^{13} \mathrm{CO}$ is a tracer of moderately dense $\left(\sim 10^{3} \mathrm{~cm}^{-3}\right)$ gas. It is generally optically thin and hence is a good tracer of column density. It probes the outer regions of dense cores ( 10 to 20 K ), but not the densest interiors, where a small percentage of the total ${ }^{13} \mathrm{CO}$ mass is depleted within a small percentage of the total size $\left(\sim 10^{3} \mathrm{AU}\right.$; Tafalla et al. 2002). The most important submillimeter measurement is the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ transition: this line is the lowest-lying of the mid $-J C O$ lines that constrain density rather strongly. These data have been used to determine meaningful conclusions about the physical properties of the molecular cores. The detection of the ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ transition was unlikely, since the molecular cores are




Right Ascension (J2000)


Fig. 4.-Maps in two transitions of sources chosen for LVG analysis: ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ (left panels) and ${ }^{12} \mathrm{CO}(4 \rightarrow 3)($ right panels). Contours are in steps of $10 \%$ of the peak intensity, listed in each map. The effective beamsizes are indicated by the circles and are $216^{\prime \prime}$ in ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ and $117^{\prime \prime}$ in ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$. Maps are provided for DC 259.5-16.4, DC 274.2-0.4, DC 339.1+11.7, L100, CG 31A-C, CG 30, DC 267.5-7.4, DC 267.4-7.5, and B59.


Fig. 4-Continued


Fig. 4-Continued


Fig. 4-Continued


Fig. 4-Continued
rather cold. From our data, we were able to put upper limits on the intensities in this transition, which were important for a LVG analysis.

We find typical ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ line widths to be $1.8-2.3 \mathrm{~km} \mathrm{~s}^{-1}$ for cores with embedded stars and $1.6-1.9 \mathrm{~km} \mathrm{~s}^{-1}$ for all others. The slightly higher values for sources with embedded stars result from warmer gas whose emission is more subject to thermal and turbulent broadening of the lines. Only 10 out of 40 cores have line widths outside the $1.5-2.1 \mathrm{~km} \mathrm{~s}^{-1}$ window, the smallest being $0.8 \mathrm{~km} \mathrm{~s}^{-1}$ and the largest being $2.7 \mathrm{~km} \mathrm{~s}^{-1}$. The largest line width is found in DC $275.9+1.9$, a core with an embedded star; the smallest is found in the rather weak source DC 253.8-10.9, with no signs of star formation.

The distribution of the ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ line width is much broader. Typical values lie in the range of $1.6-2.3 \mathrm{~km} \mathrm{~s}^{-1}$ for cores with no embedded sources, and $2.6-3.7 \mathrm{~km} \mathrm{~s}^{-1}$ for cores with embedded stars. The distribution of their line widths is rather flat, with the largest line width of $5.8 \mathrm{~km} \mathrm{~s}^{-1}$ occurring in DC $303.8-14.2$. From the ${ }^{13} \mathrm{CO}$ line intensities, we were able to derive the ${ }^{13} \mathrm{CO}$
column density and the $\mathrm{H}_{2}$ column density (§ 3 ). The mean of the ${ }^{13} \mathrm{CO}$ column density is found to be $1.9 \times 10^{15} \mathrm{~cm}^{-2}$ for cores with embedded stars and $1.5 \times 10^{15} \mathrm{~cm}^{-2}$ for all others. Due to its exceptionally strong ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$ line intensity, B59 also stands out here with a ${ }^{13} \mathrm{CO}$ column density of $4.3 \times 10^{15} \mathrm{~cm}^{-2}$ and an $\mathrm{H}_{2}$ column density of $7.3 \times 10^{21} \mathrm{~cm}^{-2}$.

For nine cores, we determined the kinetic temperature and molecular hydrogen density through an LVG analysis. For the ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ transition, we determined upper intensity limits, since we have no significant detection of this line. As a result, all temperatures and densities are also upper limits. In addition, we give a lower temperature limit, assuming the absence of ${ }^{12} \mathrm{CO}(7 \rightarrow 6)$ emission. For these nine cores, four of which have embedded stars, we find an average lower temperature limit of 16 K and an average upper limit of 26 K . The lower temperature limit is smaller than 21 K for all cores. The average upper limit on the molecular hydrogen density was found to be $10^{3.3} \mathrm{~cm}^{-3}$ for all cores. The average lower limit on the molecular hydrogen density was found to be $10^{2.9} \mathrm{~cm}^{-3}$. Due to the small sample size, we cannot detect a

TABLE 4
Results from the LVG Model

| Source | $\alpha(\mathrm{J} 2000.0)^{\mathrm{a}}$ | $\delta(\mathrm{J} 2000.0)^{\mathrm{a}}$ | $\begin{aligned} & T_{2 \rightarrow 1}^{13} \\ & (\mathrm{~K}) \end{aligned}$ | $T_{4 \rightarrow 3}^{12}$ <br> (K) | $\begin{gathered} T_{7 \rightarrow 6}^{12} \\ (\mathrm{~K}) \end{gathered}$ | $T_{7 \rightarrow 6}^{12} / T_{4 \rightarrow 3}^{12}$ <br> (K) | $T_{4 \rightarrow 3}^{12} / T_{2 \rightarrow 1}^{13}$ <br> (K) | $\begin{aligned} & T_{\text {kin }} \\ & (\mathrm{K}) \end{aligned}$ | $T_{\text {kin }}$ <br> (K) | $\begin{gathered} \log n\left(\mathrm{H}_{2}\right) \\ \left(\mathrm{cm}^{-3}\right) \end{gathered}$ | $\begin{gathered} \log n\left(\mathrm{H}_{2}\right) \\ \left(\mathrm{cm}^{-3}\right) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  | u.1. ${ }^{\text {b }}$ | u.1. ${ }^{\text {b }}$ |  | u.1. ${ }^{\text {c }}$ | 1.1. ${ }^{\text {b }}$ | u.1. ${ }^{\text {b }}$ | 1.1. ${ }^{\text {b }}$ |
| DC 259.5-16.4........... | 071403.2 | -48 2740 | 1.6 | 0.8 | 0.5 | 0.6 | 0.5 | 22 | 14 | 3.2 | 2.9 |
| CG 31A-C................. | 080849.2 | -35 5751 | 3.1 | 4.3 | 1.4 | 0.3 | 1.4 | 28 | 17 | 3.2 | 2.8 |
| CG 30*c...................... | 080926.3 | -36 0351 | 1.6 | 1.8 | 0.8 | 0.4 | 1.2 | 30 | 18 | 3.3 | 2.9 |
| DC 267.4-7.5*........... | 082548.3 | -50 5948 | 2.5 | 2.8 | 0.3 | 0.1 | 1.1 | 22 | 14 | 3.2 | 2.9 |
| DC 267.5-7.4............. | 082646.5 | -5104 01 | 1.6 | 1.5 | 0.4 | 0.3 | 0.9 | 26 | 15 | 3.2 | 2.9 |
| DC 274.2-0.4............. | 092847.0 | $-513510$ | 1.0 | 0.6 | 0.2 | 0.4 | 0.6 | 18 | 12 | 3.7 | 3.0 |
| DC 339.1+11.7 ........... | 155911.2 | -37 3657 | 2.3 | 3.4 | 1.3 | 0.4 | 1.5 | 34 | 19 | 3.1 | 2.8 |
| B59* .......................... | 171119.8 | -27 2558 | 4.0 | 4.1 | 1.9 | 0.5 | 1.0 | 28 | 16 | 3.3 | 2.9 |
| L100* ........................ | 171607.4 | -20 5721 | 1.7 | 2.3 | 0.6 | 0.3 | 1.4 | 30 | 20 | 3.1 | 2.7 |

[^4]

Fig. 5.-Maps of DC $253.6+2.0$, DC $297.7-2.8$, DC $302.6-15.9$, and DC $303.8-14.2$ in ${ }^{13} \mathrm{CO}(2 \rightarrow 1)$. The effective beam size of $216^{\prime \prime}$ is indicated by the circle.

TABLE 5
Masses Derived from the $\mathrm{H}_{2}$ Column Densities

| Source | $\begin{gathered} \text { Distance }^{\mathrm{a}} \\ \quad(\mathrm{pc}) \end{gathered}$ | Radius ${ }_{1}{ }^{\text {b }}$ (pc) | $\begin{gathered} \text { Radius }_{2}{ }^{\mathrm{b}} \\ (\mathrm{pc}) \end{gathered}$ | $\begin{gathered} N\left(\mathrm{H}_{2}\right) \\ \left(10^{21} \mathrm{~cm}^{-2}\right) \end{gathered}$ | Mass $\left(M_{\odot}\right)$ | $\begin{gathered} T_{\text {kin }}(\mathrm{u} .1 .)^{\mathrm{c}} \\ (\mathrm{~K}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| DC 259.5-16.4.................... | $450 \pm 50$ | 0.9 | 0.7 | 1.3 | $48 \pm 12$ | 22 |
| CG $30 *$ * | $450 \pm 50$ | 0.8 | 0.7 | 1.8 | $58 \pm 14$ | 30 |
| DC 267.4-7.5* .................... | $450 \pm 50$ | 1.4 | 1.2 | 2.9 | $255 \pm 60$ | 22 |
| DC 253.6+2.0 ..................... | $450 \pm 50$ | 0.9 | 0.7 | 2.4 | $103 \pm 26$ | ... |
| DC 274.2-0.4...................... | $450 \pm 50$ | 0.8 | 0.7 | 1.5 | $56 \pm 14$ | 18 |
| DC 297.7-2.8..................... | $150 \pm 30$ | 0.2 | 0.3 | 3.3 | $12 \pm 4$ | ... |
| DC 302.6-15.9.................... | $150 \pm 30$ | 0.2 | 0.3 | 0.9 | $4 \pm 2$ | $\ldots$ |
| DC 303.8-14.2* .................. | $178 \pm 18$ | 0.5 | 0.4 | 1.5 | $19 \pm 4$ | . |
| L100* ................................. | $125 \pm 25$ | 0.4 | 0.4 | 1.9 | $14 \pm 6$ | 30 |

[^5]significant difference between the mean densities of cores with and without embedded stars.

We have obtained masses for nine cores. They range from 4 to $255 M_{\odot}$, depending on the size of the core. A correlation between temperature or density and mass cannot be established, which is mainly due to the small sample size.

In general, the sources in this survey behave rather uniformly with only very few exceptions. Cores with embedded sources tend to have larger line widths and higher ${ }^{13} \mathrm{CO}$ column density than cores with no embedded stars.

From the comparison to previously published CO surveys of molecular cores (Keto \& Myers 1986; Vilas-Boas et al. 1994, 2000), we can say that the c2d sample of molecular cores has a range of properties (line width, column density, size, mass, embedded stars) similar to those of past studies. Only three sources from the sample covered in this survey have previously been observed in other CO surveys, albeit in lower rotational transitions. To obtain even more accurate results with an LVG analysis, a sur-
vey of these sources at one other CO transition would be very valuable.

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[^6]
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[^1]:    ${ }^{\text {a }}$ Coordinates of the map center. Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.
    ${ }^{\mathrm{b}}$ Cores with known embedded sources as detected using IRAS are marked with asterisks ( ${ }^{*}$ ).
    ${ }^{\text {c }}$ Data for ${ }^{12} \mathrm{CO}(4 \rightarrow 3)$ on DC $297.7-2.8$ will be presented by T. L. Bourke et al. (2007, in preparation).

[^2]:    ${ }^{3}$ See http://arcsec.sejong.ac.kr/~skim/lvg2.html.

[^3]:    ${ }^{\text {a }}$ Coordinates of spectrum with highest integrated intensity. Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.
    ${ }^{\mathrm{b}}$ Line width obtained using high-resolution spectrometer data.
    ${ }^{\text {c }}$ Cores with known embedded sources as detected using IRAS are marked with asterisks $\left(^{*}\right)$.

[^4]:    ${ }^{\text {a }}$ Coordinates of position at which LVG analysis was performed. Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.
    b Upper and lower limits are denoted "u.l." and "1.1."
    ${ }^{\text {c }}$ Cores with known embedded sources as detected using IRAS are marked with asterisks $\left(^{*}\right)$.

[^5]:    ${ }^{\text {a }}$ Distances as derived in Woermann et al. (2001), Franco (1995), Corradi et al. (1997), Whittet et al. (1997), Knude \& Høg (1998), and de Geus et al. (1989).
    ${ }^{\mathrm{b}}$ Radii of the elliptical half power width of the core.
    ${ }^{\text {c }}$ Upper limit as derived from the LVG analysis.
    ${ }^{\mathrm{d}}$ Cores with known embedded sources as detected using IRAS are marked with asterisks ( ${ }^{*}$ ).

[^6]:    ${ }^{4}$ See http://skyview.gsfc.nasa.gov.

