

1. Basic Atmospheric Properties

Cf. **Houghton**, Chapter 1; **Goody & Yung**, Chapter 1; **Brasseur and Solomon**, Chapter 3 (Structure and Dynamics in Edition 2; Dynamics and Transport in Edition 3).

Hydrostatic equilibrium

The Earth's atmosphere is, and other planetary atmospheres are, normally well-described as being close to a state of *hydrostatic equilibrium*, wherein the relationship of pressure, P , altitude, z , and temperature, T , is given by

$$dP = -g\rho dz, \quad \rho = \frac{\bar{m}P}{RT}. \tag{0.1}$$

Then

$$\frac{dP}{P} = d \ln P = -\frac{dzg\bar{m}}{RT} \equiv -\frac{dz}{H}. \tag{0.2}$$

g is the gravitational acceleration corresponding to the altitude (it varies little over the distance corresponding to the bulk of the atmosphere, except for the gas giant planets), ρ the gas density, \bar{m} the mean molecular weight of the gas (28.9644 grams per mole, for normal air), R the ideal gas constant (82.057361 cm³ atmosphere mole⁻¹ °K⁻¹), and H is the local pressure *scale height*, defined to describe the change of pressure with altitude. A typical range of scale heights in the lower and middle Earth's atmosphere is 7-8 km.

Thus, hydrostatic equilibrium can be conveniently used to relate pressure and temperature (or to determine the overall atmospheric structure, given two of the three variables z , P , and T). Density scale heights may also be defined. Scale heights are sometimes used to describe variations of the vertical profiles for different gases, which can vary with gas, sometimes dramatically, as well as with height. In this case, they are not determined by the above simple thermodynamic relationship. *Note the use of the ideal gas law to determine ρ . How much would the use of the van der Waals correction for non-ideal behavior change the P,T relationship in the Earth's atmosphere?*

Adiabatic lapse rate

The decrease in atmospheric temperature from purely thermodynamic considerations can be derived by considering the work done by air as it expands under conditions where the heat content, q , does not change, $dq = 0$, that is, adiabatic expansion.

$$\frac{dT}{dz} = \frac{-g}{c_p} \approx -10^\circ K / km, \tag{0.3}$$

where c_p is the heat capacity at constant pressure (See **Houghton 1.4**, for the full derivation).

Basic structure and variability

Figure 1.1 shows the approximate shape and range of variability of the height (and pressure) versus temperature profile for the Earth's atmosphere. Note particular the higher, and colder, tropopause at summer and low latitudes versus the lower, and warmer, tropopause at high

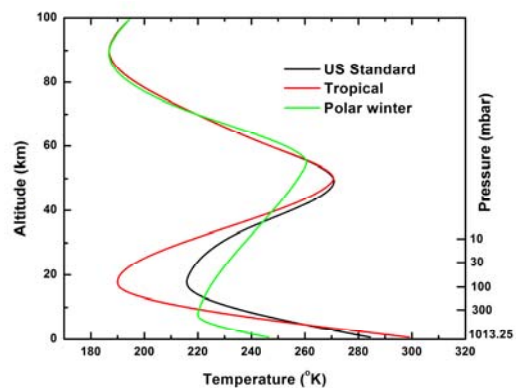


Figure 1.1 Basic structure and variability

rows are megahertz, gigahertz, and terahertz. The distinction between the use of frequency versus wavelength is often due to instrument technique and/or spectral range: radiofrequency (RF), microwave, heterodyne techniques \Rightarrow frequency; Fourier transform spectrometer (FTS) \Rightarrow frequency (usually cm^{-1}), since their spectra are linear with energy. Dispersive instruments (grating/prism spectrometers) \Rightarrow wavelength, since their spectra are linear with wavelength (approximately so in the case of prisms, more exactly so for gratings).

To emphasize, the relationships $E = h\nu = hc\sigma = hc/\lambda$ tell how to *convert among units*. Thus, a frequency of 200 cm^{-1} corresponds to a frequency of 5.995849×10^{12} Hz (5.995849 THz, 5.995849×10^3 GHz, 5.995849×10^6 MHz), and a wavelength of 0.005 cm ($50\text{ }\mu\text{m}$, $50,000$ nm, $500,000\text{ }\text{\AA}$).

Homework problem (assigned January 26, due February 7) on units – MKS versus cgs, wavelengths versus wavenumbers and frequencies.

Construct a table showing wavelengths and frequencies (nm, μm , cm^{-1} , MHz, GHz, and THz) for: CO $1 \rightarrow 0$ and $2 \leftarrow 0$ band centers (2143.272 cm^{-1} ; 4260.063 cm^{-1}); ClO MLS emission line (204.35 GHz); O₂ A band center (13120.909 cm^{-1}); CO₂ $15\text{ }\mu\text{m}$ “greenhouse” band (667.380 cm^{-1}); O₃ TOMS “on” wavelength (317.35 nm).