# AN IRAS HIGH RESOLUTION IMAGE RESTORATION (HIRES) ATLAS OF ALL INTERACTING GALAXIES IN THE IRAS REVISED BRIGHT GALAXY SAMPLE 

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#### Abstract

The importance of far-infrared observations for our understanding of extreme activity in interacting and merging galaxies has been illustrated by many studies. Even though two decades have passed since its launch, the most complete all-sky survey to date from which far-IR selected galaxy samples can be chosen is still that of the Infrared Astronomical Satellite (IRAS). However, the spatial resolution of the IRAS all-sky survey is insufficient to resolve the emission from individual galaxies in most interacting galaxy pairs, and hence previous studies of their far-IR properties have had to concentrate either on global system properties or on the properties of very widely separated and weakly interacting pairs. Using the HIRES image reconstruction technique, it is possible to achieve a spatial resolution ranging from $30^{\prime \prime}$ to 1.5 (depending on wavelength and detector coverage), which is a fourfold improvement over the normal resolution of $I R A S$. This is sufficient to resolve the farIR emission from the individual galaxies in many interacting systems detected by $I R A S$, which is very important for meaningful comparisons with single, isolated galaxies. We present high-resolution $12,25,60$, and $100 \mu \mathrm{~m}$ images of 106 interacting galaxy systems contained in the IRAS Revised Bright Galaxy Sample (RBGS, Sanders et al.), a complete sample of all galaxies having a $60 \mu \mathrm{~m}$ flux density greater than 5.24 Jy . These systems were selected to have at least two distinguishable galaxies separated by less than three average galactic diameters, and thus we have excluded very widely separated systems and very advanced mergers. Additionally, some systems have been included that are more than three galactic diameters apart, yet have separations less than $4^{\prime}$ and are thus likely to suffer from confusion in the RBGS. The new complete survey has the same properties as the prototype survey of Surace et al. We find no increased tendency for infrared-bright galaxies to be associated with other infrared-bright galaxies among the widely separated pairs studied here. We find small enhancements in far-IR activity in multiple galaxy systems relative to RBGS noninteracting galaxies with the same blue luminosity distribution. We also find no differences in infrared activity (as measured by infrared color and luminosity) between late- and early-type spiral galaxies.


Key words: atlases - galaxies: interactions - infrared: galaxies - infrared: general techniques: image processing
On-line material: machine-readable table

## 1. INTRODUCTION

In the last two decades it has become apparent that interactions between galaxies can play a significant role in their evolution. From the early dynamical simulations of Toomre \& Toomre (1972) to more modern work by Barnes, Hernquist, and others (Barnes \& Hernquist 1992, and references therein), it has become apparent that interactions and mergers between galaxies can radically alter their morphology by inducing shells, bars, tails, and other tidal features. Perhaps more importantly, cancellation of angular momentum during the merger process can lead to a radical redistribution of the gas content of the galaxies, with very rapid gas inflow into the galaxy cores. This supply of fresh material could possibly fuel an active galactic nucleus or provide the high gas densities needed to lead to a sudden burst of star formation.

There is considerable evidence that enhanced star formation is associated with interacting galaxies (Sulentic 1988, and
references therein). The young OB stars that dominate the starburst radiate primarily in the optical and ultraviolet, but surrounding gas and dust reprocesses this radiation and thus strongly radiates at thermal wavelengths in the far-infrared. Far-IR luminosity is thus indicative of the magnitude of recent star formation activity (Telesco 1988; Lonsdale et al. 1984). In addition, because of the increased temperature of the heated dust, we expect the far-IR colors to be a good diagnostic of enhanced star formation. Therefore, many studies have concentrated on the far-IR properties of interacting galaxies.

Several studies have also discussed the incidence of multiple bright galaxies being found within a given interacting galaxy system. Haynes \& Herter (1988) found that for galaxy pairs separated by $2^{\prime}-10^{\prime}$, approximately $10 \%$ have multiple components brighter than 0.5 Jy at $60 \mu \mathrm{~m}$ and 1 Jy at $100 \mu \mathrm{~m}$. Xu \& Sulentic (1991) also concluded that in the majority of interacting systems, only one galaxy is infrared bright. These results agree with earlier work by Joseph et al. (1984), which
concluded, based on near-IR colors, that usually only one galaxy in a pair showed signs of unusual activity. This is an interesting result, because it suggests that specific properties of the interacting galaxies may determine whether or not they become emitters in the far-IR as well. Testing this hypothesis requires resolution of the individual galaxies in the far-IR, which is the goal of this $I R A S$ study.

The canonical figure used by many authors to delineate interacting versus noninteracting systems is a projected separation of three average galactic diameters, as presumably galaxies this close to one another are also close enough to exert a considerable gravitational effect (Dahari 1984; Byrd et al. 1987; Surace et al. 1993). However, for most galaxies detected by $\operatorname{IRAS}$, this typically corresponds to an angular separation of a few arcminutes, which is less than the resolution normally achieved by $I R A S$ using the one-dimensional coadder ADDSCAN or the two-dimensional FRESCO imaging process. As a result, it has been impossible to study the far-IR properties of the individual galaxies, and most studies have either made assumptions about the distribution of flux between galaxies within the interacting system (Bushouse 1986) or have concentrated on widely separated pairs (Haynes \& Herter 1988; Xu \& Sulentic 1991). Since previous studies of very widely separated galaxy pairs indicate that in a substantial fraction of interacting systems only one galaxy is unusually active in the far-IR (Xu \& Sulentic 1991), it is necessary to resolve these galaxies in order to properly study those properties such as morphology that are unique to the individual galaxies. Additionally, Xu \& Sulentic (1991) found evidence that at smaller separations (and hence greater interaction strengths), there was a greater enhancement of far-IR activity. Therefore it would be valuable if these studies could be extended to smaller separations for which more observable changes are taking place.

Development of the maximum correlation method algorithm (MCM; Aumann et al. 1990) for use in IRAS image reconstruction significantly increased the resolution of IRAS observations. As implemented in the HIRES process, MCM is an iterative image reconstruction technique that involves using the known response functions of the $\operatorname{IRAS}$ detectors to scan simulated image estimates, which are then compared to the actual detector data. In this way, a high-resolution image estimate is formed. The result is typically a fivefold increase in resolution, varying from roughly $30^{\prime \prime} \times 45^{\prime \prime}$ at $12 \mu \mathrm{~m}$ to $72^{\prime \prime} \times 130^{\prime \prime}$ at $100 \mu \mathrm{~m}$, with the actual achieved resolution being highly dependent on the geometry of the detector coverage (Surace et al. 1993). Unfortunately, the HIRES process is extremely computer intensive. When developed, a single field typically took a day or more to process. As a result, the earlier work by Surace et al. (1993) was rather limited in scope, with only 23 systems being resolvable. On a modern computer, this computing time is reduced to approximately 15 minutes, thus making feasible processing of a substantially larger sample. ${ }^{1}$

In $\S 2$ we present the sample selection criteria for the objects examined here and the data reduction techniques used for reconstructing the IRAS images and measuring the galaxy fluxes. The fluxes at each $\operatorname{IRAS}$ wavelength are presented in tabular form, and contours of the infrared emission are shown overlaid on optical images of the galaxies. In § 3 we present

[^0]properties of the catalog and some results derived from them. Appendix A presents additional notes for selected galaxy systems. Finally, in Appendix B we include data for galaxy systems that were originally included in the Bright Galaxy Sample (BGS) but were subsequently dropped from the Revised Bright Galaxy Sample (RBGS) after a reanalysis of their fluxes. These objects are provided for the interest of the reader but do not bear on the analysis of the catalog.

## 2. DATA

### 2.1. Sample

All of the targets were selected from the $I R A S$ RBGS (Sanders et al. 2003). The RBGS consists of all 629 galaxies detected by $I R A S$ with a $60 \mu \mathrm{~m}$ flux density greater than 5.24 Jy and thus is similar to and includes all of the well-studied BGS (Soifer et al. 1989), but extends coverage to the entire sky at Galactic latitudes $|b|>5^{\circ}$.

The following criterion was applied in order to select close pairs from the RBGS:

$$
\begin{equation*}
\frac{2 S_{12}}{D_{1}+D_{2}} \leq 3 \tag{1}
\end{equation*}
$$

where $S_{12}$ is the distance between galaxy centers and $D_{1}$ and $D_{2}$ are their optical diameters, as measured from the Palomar Sky Survey. This criterion therefore selects all systems in which the galaxies are separated by less than three times their average diameter. Note that this excludes very advanced mergers such as Arp 220 in which the individual galactic disks can no longer be distinguished. This also has the additional benefit of selecting systems that are sufficiently separated as to be resolvable with HIRES. As such, the sample includes all of the galaxies listed in Table 1 of Surace et al. (1993). Additionally, in an attempt to resolve sources listed in the RBGS that were likely to be confused because of small separations, we included all small galaxy pairs with apparent separation less than $4^{\prime}$. This separation was determined by the normal survey resolution of $4^{\prime}$, which in turn is set by the IRAS $100 \mu \mathrm{~m}$ detector size.

### 2.2. Data Reduction

The IRAS data were processed in a manner similar to Surace et al. (1993). The raw detector scans were initially extracted from the $I R A S$ database using the SNIPSCAN process. These raw detector scans were then flattened using an iterative fitting technique that removed the detector baselines, and they were then deglitched in order to remove artifacts such as cosmic ray hits using the LAUNDR process. The HIRES process was then applied to the detector scans. Restoration was done on $1 \mathrm{deg}^{2}$ fields in order to improve detector baseline coverage, with a pixel size of $15^{\prime \prime}$ per pixel, which is sufficient to adequately sample the restored $I R A S$ beam. The algorithm was iterated 20 times as further iterations tend only to increase noise amplification with little improvement in resolution.

In order to aid in the interpretation of the IRAS data, optical images were extracted from the Digital Sky Survey (DSS), and the $\operatorname{IRAS}$ data were overlaid on them (Fig. 1). This was valuable in interpreting the correspondence between the resolved IRAS objects and the optical galaxies. In some cases there were small, uncataloged optical galaxies in the DSS images, and the DSS images were used to derive their positions. The optical images have a pixel size of 1.7 . The astrometry of the optical images is based on a linear approximation to the polynomial


FIg. 1.-HIRES data for each pair or group in the RBGS overlaid on gray-scale images from the Palomar Digital Sky Survey. The contours are at $10 \%$, $16 \%$, $25 \%, 40 \%$, and $63 \%$ of the peak flux; they show the results of 20 iterations of the MCM algorithm and represent the highest achievable IRAS resolution at 12 , 25 , 60 , and $100 \mu \mathrm{~m}$. The axis labels are B1950 coordinates. The scale bar inside each $100 \mu \mathrm{~m}$ panel represents 50 or 25 kpc at the distance of the system (as labeled).



Fig. 1.-Continued


Fig. 1.-Continued


Fig. 1.-Continued




Fig. 1.-Continued



Fig. 1.-Continued


Fig. 1.-Continued



Fig. 1.-Continued



Fig. 1.-Continued


$12 \mu \mathrm{~m}$




Fig. 1.-Continued



Fig. 1.-Continued


Fig. 1.-Continued

plate solution provided by the Space Telescope Science Institute (STScI) and yields positions accurate to roughly $2^{\prime \prime}$ (Laidler et al. 1994). The astrometry of the IRAS images is limited by the pointing accuracy of the satellite, which was approximately $2^{\prime \prime}$ in the in-scan direction and $10^{\prime \prime}$ cross-scan (Beichmann et al. 1988). Furthermore, the astrometry of pointlike sources produced by the HIRES technique is known observationally to be approximately $20^{\prime \prime}$ (Laughlin et al. 1990). Thus, registration of the images should be accurate to within one or two HIRES pixel elements and certainly should be better than the typical IRAS beam size. Optical identifications were made using the coordinates and names given in the NASA/IPAC Extragalactic Database (NED), which are in turn derived from the Third Reference Catalogue of Bright Galaxies (RC3, de Vaucouleurs et al. 1991). When no identifications were available, the galaxies were identified directly from the DSS images and are labeled from northeast to southwest.

Photometry was accomplished in two ways. When the galaxies were cleanly separated, aperture photometry was performed via the IPAC Skyview software using polygonal apertures of a size sufficient to insure that all of the galaxy flux was measured. In those cases for which components appeared to be resolved but not separated, the data were modeled with two-dimensional elliptical Gaussians using the AIPS IMFIT and JMFIT routines. This is justified in that although the geometry of the $I R A S$ beam is variable, it usually has roughly the form of an elliptical Gaussian whose exact size and orientation depend on the detector scan geometry. Peak positions were constrained to the position indicated by the nearest separated $\operatorname{IRAS}$ wavelength. If none of the $\operatorname{IRAS}$ data were able to supply positions, then the Gaussian centers were constrained to the locations of the optical peaks as given by the RC3, when possible, and otherwise according to positions measured directly from the DSS. Note that the latter could introduce a bias in that it presupposes correspondence between the infrared and optical centers. In cases in which the galaxies are well separated at optical wavelengths, the optical and infrared peaks do correspond. Many of the galaxies that were decomposed using Gaussian fitting are also well separated at optical wavelengths but are too close to one another to be separated by $\operatorname{IRAS}$. In these cases it is reasonable to assume that the optical peaks correspond to the infrared peaks. Only in cases of advanced mergers, such as NGC 4038, would this assumption break down.

Table 1 presents the measured global photometry for each of the galaxies identified in the $\operatorname{IRAS}$ images from the RBGS sample. The supporting data in Table 1 and Table 2 were taken
from the NED, and relevant notes regarding this database are given below. It should be noted that the magnitudes, morphological types, etc., listed by NED are generally not on any homogeneous system, although when possible, data from NED are derived from the RC3. The table is ordered by increasing B1950 right ascension of the galaxy systems, as given by the westernmost component. The columns are as follows:

Column (1).-The galaxy name. Names are given in order of preference from the NGC, UGC, Catalog of Galaxies and Clusters of Galaxies (CGCG), Morphological Catalogue of Galaxies (MCG), Markarian (Mrk), and the Two Micron All Sky Survey (2MASS) catalogs. Relevant cross-identifications are also given. The given names are those associated with a specific coordinate as given by the RC3, NED, or the Arp atlas. The Arp name associated with a given galaxy group is listed with the first (westernmost) object, but no special significance is indicated by this.

Columns (2) and (3).-Equinox 1950 coordinates. The given coordinates are those of the centroid of the infrared emission in cases in which there was a separated infrared detection. Otherwise, the optical position from NED is listed. In almost all cases these coordinates originate from the RC3, although currently NED lists values from 2MASS. If NED listed no coordinates, then the coordinates are those measured directly from the DSS. Equinox B1950 coordinates were chosen because this is the epoch of the positional calibration for the $\operatorname{IRAS}$ Level 1 archive scans used to construct the HIRES images.

Column (4).-The position type. If "O," then the given position is an optical center, and if "I," then a new infrared center is derived from the HIRES images.

Column (5).-Radial velocity in $\mathrm{km} \mathrm{s}^{-1}$. In all cases in which these are listed, they are spectroscopic redshifts from a variety of sources as given by NED. Uncertainties are typically $10-100 \mathrm{~km} \mathrm{~s}^{-1}$.

Column (6).-Total optical magnitude as given by NED. In most cases this is the blue magnitude listed in the RC3.

Column (7)-(14).-12, 25, 60, and $100 \mu \mathrm{~m}$ integrated fluxes in Jy and the associated $1 \sigma$ uncertainties. In cases in which the targets were also given by Surace et al. (1993), the objects were remeasured in order to ensure uniformity of calibration with the rest of the survey (see $\S 2.3$ ). In some cases the galaxies were still unresolved at all wavelengths, so only the global flux as measured from the HIRES data is given. In cases in which one or more components remained unresolved but some components were resolved, the brightest unresolved component in the resolved waveband gives the flux of all the components, and subsequent unresolved components are marked with ellipses.

Integrated Flux Densities of RBGS Interacting Galaxies

| Name <br> (1) | $\begin{aligned} & \text { R.A. } \\ & \text { (B1950) } \end{aligned}$ (2) | $\begin{aligned} & \text { Decl. } \\ & \text { (B1950) } \\ & (3) \end{aligned}$ | Position Type <br> (4) | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ <br> (5) | $m_{z}$ <br> (6) | $\begin{gathered} 12 \mu \mathrm{~m} \\ (\mathrm{Jy}) \\ \text { (7) } \end{gathered}$ | $\begin{gathered} \sigma_{12} \\ (\mathrm{Jy}) \\ (8) \end{gathered}$ | $\begin{gathered} 25 \mu \mathrm{~m} \\ (\mathrm{Jy}) \\ \text { (9) } \end{gathered}$ | $\sigma_{25}$ <br> (Jy) <br> (10) | $\begin{gathered} 60 \mu \mathrm{~m} \\ (\mathrm{Jy}) \\ (11) \end{gathered}$ | $\begin{gathered} \sigma_{60} \\ (\mathrm{Jy}) \\ (12) \end{gathered}$ | $100 \mu \mathrm{~m}$ (Jy) <br> (13) | $\sigma_{100}$ <br> (Jy) <br> (14) | $L_{\text {FIR }}$ <br> $\left(L_{\odot}\right)$ <br> (15) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| MCG -02-01-052 = VV 352 | 001617.2 | -1038 21 | I | 8193 | 13.6 | $<0.24$ |  | 0.15 | 0.08 | $<0.25$ | $\ldots$ | $<0.50$ |  |  |
| MCG -02-01-051 | 001618.0 | -10 3916 | I | 8112 | 14.8 | $<0.24$ | . | 1.19 | 0.06 | 7.35 | 0.12 | 10.22 | 1.35 | 11.13 |
| NGC 230. | 003959.0 | -23 5412 | I | 6250 | 15.4 | $<0.30$ |  | 0.38 | 0.08 | 1.15 | 0.09 | 3.10 | 0.42 | 10.22 |
| NGC 232. | 004017.0 | -23 5006 | I | 6047 | 14.5 | 0.32 | 0.08 | 1.22 | 0.10 | 9.34 | 1.80 | 15.75 | 4.80 | 11.01 |
| NGC 235. | 004024.0 | -23 4854 | I | 6664 |  | $<0.18$ |  | 0.55 | 0.10 | 2.67 | 0.52 | 3.98 | 1.24 | 10.53 |
| MCG +12-02-001 | 005040.7 | 724847 | I | 4706 | 16.0 | 0.84 | 0.10 | 3.84 | 0.15 | 22.93 | 0.25 | 31.75 | 3.00 | 11.15 |
| NGC 317a = UGC 593 | 005449.8 | 433151 | O | 5293 | 15.0 | $<0.20$ |  | $<0.25$ |  | ... | $\ldots$ |  | $\ldots$ |  |
| NGC 317b $=$ UGC 594 | 005451.1 | 433123 | I | 5334 | 13.8 | 0.28 | 0.04 | 1.15 | 0.15 | 9.34 | 1.00 | 13.95 | 1.40 | 10.88 |
| CGCG 536-014. | 005452.9 | 432547 | I | 5413 | 15.1 | $<0.20$ |  | $<0.27$ |  | 0.78 | 0.90 | 1.97 | 0.20 | 9.91 |
| IC $1623 \mathrm{a} / \mathrm{b}$. | 010521.0 | -174623 | I | 6016 | 14.3 | 0.92 | 0.10 | 3.86 | 0.20 | 23.85 | 0.20 | 31.53 | 0.25 | 11.37 |
| IC 1622 | 010509.0 | -174818 | O | 6343 | 15.1 | $<0.30$ |  | $<0.20$ |  | 0.50 | 0.20 | 1.25 | 0.25 | 9.86 |
| NGC $520^{\text {a }}$ | 012159.6 | 033151 | I | 2281 | 12.0 | 1.07 | 0.20 | 3.08 | 0.45 | 31.62 | 2.00 | 47.76 | 3.30 | 10.68 |
| NGC 633. | 013410.6 | -3734 33 | I | 5148 | 13.5 | 0.22 | 0.03 | 0.72 | 0.08 | $2.01{ }^{\text {b }}$ | 0.54 | $2.32{ }^{\text {b }}$ | 1.00 | 10.15 |
| ESO 297-G012 | 013412.3 | -37 3541 | I | 5193 | 14.9 | 0.34 | 0.02 | 0.96 | 0.07 | $6.24{ }^{\text {b }}$ | 1.40 | $11.01^{\text {b }}$ | 4.20 | 10.71 |
| UGC 1385 = Mrk 2. | 015156.3 | 364018 | I | 5621 | 13.9 | 0.39 | 0.06 | 1.15 | 0.14 | 5.92 | 0.70 | 9.27 | 1.10 | 10.74 |
| CGCG 522-061. | 015159.1 | 363708 | I | 4919 | 15.2 | $<0.28$ | $\ldots$ | $<0.20$ | ... | 0.51 | 0.07 | $<0.76$ | $\ldots$ | ... |
| CGCG 522-062. | 015204.3 | 364031 | O | 5621 | 15.2 | $<0.20$ | ... | $<0.25$ | . | $<0.25$ | $\ldots$ | $<0.85$ | $\ldots$ | $\ldots$ |
| NGC $833=$ Arp 318 | 020653.3 | -1022 09 | O | 3864 | 13.7 | $<0.30$ | $\ldots$ | $<0.30$ | $\ldots$ | $<0.12$ | ... | $<0.40$ | $\ldots$ |  |
| NGC $835=$ Mrk 1021 | 020657.3 | -1022 19 | I | 4152 | 12.9 | 0.45 | 0.15 | 0.97 | 0.30 | 6.33 | 0.65 | 11.54 | 1.20 | 10.53 |
| NGC 838. | 020711.1 | -10 2257 | I | 3851 | 13.6 | 0.92 | 0.15 | 2.70 | 0.30 | 12.96 | 1.40 | 16.96 | 1.50 | 10.72 |
| NGC 839. | 020715.6 | -1025 11 | I | 3847 | 13.9 | 0.68 | 0.15 | 3.34 | 0.30 | 12.80 | 1.40 | 12.90 | 1.40 | 10.68 |
| UGC 1720b. | 021118.9 | 045606 | O | 9200 | 15.5 | $<0.18$ |  | $<0.30$ |  | $<0.30$ |  | $<0.60$ |  |  |
| UGC $1720=$ IC 214. | 021129.2 | 045628 | I | 9061 | 15.4 | 0.33 | 0.10 | 0.68 | 0.23 | 5.58 | 0.25 | 9.23 | 0.30 | 11.14 |
| NGC 876. | 021510.0 | 141726 | I | 3860 | 16.5 | $\ldots$ | ... | ... |  | 4.62 | 1.30 | ... |  |  |
| NGC 877. | 021515.3 | 141901 | I | 3913 | 12.6 | 1.14 | 0.12 | 1.94 | 0.41 | 8.82 | 1.50 | 29.55 | 2.00 | 10.75 |
| MCG +05-06-035 = Mrk 1034 NED01. | 022020.8 | 315742 | O | 10083 | 15.4 | $<0.21$ |  | $<0.20$ |  | $<1.60$ | ... | ... | ... |  |
| MCG +05-06-036 = Mrk 1034 NED02. | 022023.6 | 315756 | I | 10121 | 15.0 | 0.30 | 0.08 | 0.69 | 0.11 | 6.83 | 1.10 | 12.15 | 1.20 | 11.33 |
| NGC 1097.. | 024410.6 | -30 2902 | I | 1275 | 10.2 | 4.16 | 0.42 | 9.27 | 0.95 | 58.29 | 6.00 | 114.82 | 14.00 | 10.48 |
| UGC 2369 Ned 02 | 025115.7 | 144625 | O | 9947 | 15.5 | ... | ... | $<0.40$ | $\ldots$ | <1.00 | . | ... | ... |  |
| UGC 2369 Ned 01 | 025115.9 | 144603 | I | 9804 | 15.5 | 0.37 | 0.11 | 1.62 | 0.21 | 8.27 | 0.20 | 11.75 | 0.40 | 11.35 |
| NGC 1270. | 031539.6 | 411719 | O | 4871 | 14.3 | $<0.20$ | ... | $<0.50$ | $\ldots$ | 0.38 | 0.05 | $<0.80$ | ... | ... |
| NGC 1272. | 031602.8 | 411835 | O | 4021 | 12.9 | $<0.20$ | $\ldots$ | $<0.40$ | $\ldots$ | $<0.40$ | $\ldots$ | $<0.40$ | $\ldots$ |  |
| NGC 1273. | 031608.2 | 412134 | O | 5351 | 14.3 | $<0.20$ | $\ldots$ | $<0.40$ | $\ldots$ | $<0.80$ | $\ldots$ | $<0.50$ | $\ldots$ | $\ldots$ |
| UGC 2665. | 031608.6 | 412715 | I | 7806 | 15 | $<0.20$ | $\ldots$ | $<0.40$ | $\ldots$ | 0.72 | 0.10 | 1.23 | 0.15 | 10.12 |
| IC 1907. | 031615.6 | 412358 | O | 4420 | 15.4 | $<0.20$ | $\ldots$ | $<0.40$ | $\ldots$ | $<0.40$ | $\ldots$ | $<0.50$ | $\ldots$ | ... |
| NGC 1274. | 031621.9 | 412205 | O | 6447 | 15.1 | $<0.20$ | $\ldots$ | $<0.40$ | $\ldots$ | $<0.30$ | $\ldots$ | $<0.30$ | . $\cdot$ | $\ldots$ |
| NGC 1275. | 031628.7 | 411948 | I | 5260 | 12.6 | 0.93 | 0.10 | 3.02 | 0.31 | 7.09 | 0.75 | 7.60 | 0.90 | 10.70 |
| NGC 1277/8. | 031634.8 | 412312 | I | 4982 |  | $<0.20$ |  | $<0.40$ |  | 0.52 | 0.05 | $<0.70$ | ... | ... |
| NGC 1281. | 031647.2 | 412658 | O | 4201 | 14.5 | $<0.20$ | $\ldots$ | $<0.40$ |  | $<0.90$ | $\ldots$ | $<0.90$ | $\ldots$ |  |
| NGC 1283. | 031657.0 | 411306 | O | 6749 | 14.7 | $<0.20$ |  | $<0.40$ |  | 0.43 | 0.05 | $<0.40$ | ... |  |
| IRAS 03217+4022 | 032147.6 | 402301 | I | 7007 |  | $<0.28$ |  | 1.12 | 0.17 | 7.92 | 0.37 | 13.19 | 0.41 | 11.07 |
| CGCG 468-002. | 050526.9 | 171823 | I | 5454 | 13.5 | $<0.30$ |  | 0.15 | 0.10 | 7.21 | 1.00 | 11.15 | 2.00 | 10.79 |
| UGC 3405. | 060445.3 | 802915 | I | 3791 | 15.3 | 0.14 | 0.05 | 0.20 | 0.04 | $1.93{ }^{\text {b }}$ | 0.31 | $6.08^{\text {b }}$ | 1.80 | 10.05 |
| UGC 3410. | 060519.2 | 802744 | I | 3887 | 15.0 | 0.49 | 0.05 | 0.90 | 0.07 | $7.69{ }^{\text {b }}$ | 1.15 | $18.32^{\text {b }}$ | 5.10 | 10.61 |
| NGC 2146. | 061040.1 | 782223 | I | 893 | 11.4 | 7.36 | 0.80 | 21.66 | 2.40 | 154.12 | 16.00 | 217.44 | 24.00 | 10.54 |
| NGC 2207.. | 061414.4 | -21 2115 | O | 2741 | 12.2 | 1.02 | 0.20 | 2.00 | 0.20 | 19.9 | 0.18 | 43.40 | 0.16 | 10.70 |
| IC 2163 ......... | 061420.3 | -212126 | O | 2765 | 11.6 | 0.58 | 0.20 | 1.10 | 0.20 | ... | ... | ... | ... |  |

TABLE 1—Continued

| Name <br> (1) | $\begin{aligned} & \text { R.A. } \\ & \text { (B1950) } \\ & (2) \end{aligned}$ | $\begin{aligned} & \text { Decl. } \\ & \text { (B1950) } \\ & (3) \end{aligned}$ | Position Type <br> (4) | $\begin{gathered} c z \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \\ (5) \end{gathered}$ | $m_{z}$ <br> (6) | $12 \mu \mathrm{~m}$ (Jy) <br> (7) | $\sigma_{12}$ <br> (Jy) <br> (8) | $25 \mu \mathrm{~m}$ (Jy) (9) | $\begin{gathered} \sigma_{25} \\ (\mathrm{Jy}) \\ (10) \end{gathered}$ | $60 \mu \mathrm{~m}$ (Jy) (11) | $\sigma_{60}$ <br> (Jy) <br> (12) | $100 \mu \mathrm{~m}$ (Jy) (13) | $\begin{aligned} & \sigma_{100} \\ & (\mathrm{Jy}) \\ & (14) \end{aligned}$ | $L_{\text {FIR }}$ <br> $\left(L_{\odot}\right)$ <br> (15) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NGC 2221. | 061927.5 | -57 3315 | I | 2532 | 13.8 | 0.34 | 0.05 | 0.62 | 0.20 | 5.35 | 0.50 | 11.10 | 0.50 | 10.05 |
| NGC 2222 | 061928.0 | -573034 | I | 2602 | 14.2 | 0.18 | 0.05 | 0.18 | 0.05 | 2.11 | 0.50 | 3.30 | 0.50 | 9.62 |
| ESO 161-G001 | 061945.1 | -572827 | I |  | 14.8 | $<0.15$ |  | $<0.15$ |  | 0.71 | 0.20 | 0.72 | 0.20 |  |
| ESO 255-IG007 NED 01/02/03 . | 062601.0 | -470848 | I | 11637 |  | 0.38 | 0.08 | 1.57 | 0.23 | 10.38 | 0.10 | 12.52 | 0.30 | 11.57 |
| ESO 557-G001 | 062932.3 | -173641 | I |  | 15.7 | $<0.50$ | ... | $<0.30$ |  | 1.23 | 0.20 |  |  |  |
| ESO 557-G002 | 062933.9 | -173506 | I | 6385 | 15.0 | $<0.50$ |  | 0.69 | 0.09 | 6.67 | 0.20 | 9.60 | 1.00 | 10.89 |
| NGC 2341. | 070614.4 | 204105 | I | 5227 | 13.8 | 0.66 | 0.06 | 1.13 | 0.12 | 7.62 | 1.22 | $9.88{ }^{\text {b }}$ | 2.97 | 10.75 |
| NGC 2342. | 070617.4 | 204310 | I | 5276 | 13.1 | 1.00 | 0.10 | 1.76 | 0.19 | 7.99 | 1.30 | $16.66^{\text {b }}$ | 4.40 | 10.86 |
| ESO 491-G020 | 070747.0 | -27 2918 | I | 2955 | 13.8 | 0.70 | 0.30 | 2.30 | 0.50 | 16.42 | 2.00 | 23.80 | 2.90 | 10.61 |
| ESO 491-G021 | 070749.0 | -27 2936 | O | 2847 | 13.6 | $<0.30$ |  | $<0.60$ |  |  |  |  | ... | . |
| NGC 2276. | 071018.4 | 855100 | I | 2417 | 11.9 | 1.48 | 0.20 | 2.23 | 0.20 | 14.15 | 1.23 | 31.58 | 1.73 | 10.45 |
| NCG 2300. | 071545.1 | 854831 | I | 1963 | 12.1 | $<0.36$ | $\ldots$ | $<0.30$ | $\ldots$ | $<0.30$ |  | $<0.63$ |  | ... |
| NGC 2388. | 072537.9 | 335521 | I | 4134 | 14.7 | 0.82 | 0.09 | 2.64 | 0.28 | 17.42 | 1.90 | $25.64{ }^{\text {b }}$ | 3.80 | 10.93 |
| NGC 2389. | 072548.9 | 335755 | I | 3957 | 13.4 | 0.53 | 0.07 | 0.73 | 0.10 | 3.40 | 0.35 | $5.85{ }^{\text {b }}$ | 0.88 | 10.21 |
| CGCG 058-004. | 073243.5 | 114942 | I | 4984 | 15.6 | $<0.30$ | $\ldots$ | 0.45 | 0.15 | 1.56 | 0.30 | 2.22 | 1.21 | 10.04 |
| UGC 3924. | 073251.9 | 113802 | O | 5163 | 14.7 | $<0.30$ | $\ldots$ | $<0.30$ | $\ldots$ | $<0.30$ |  | 0.30 | $\ldots$ | - |
| MCG +02-20-002 $=$ NGC 2416.. | 073255.7 | 114333 | I | 5101 | 14.1 | $<0.30$ | $\ldots$ | $<0.50$ | $\ldots$ | 1.39 | 0.15 | 2.68 | 0.50 | 10.06 |
| MCG +02-20-003 | 073257.1 | 114914 | I | 4873 | 15.0 | $<0.17$ | $\ldots$ | 0.73 | 0.15 | 8.81 | 0.30 | 12.8 | 1.40 | 10.77 |
| ESO 163-G010 | 073646.7 | -550408 | O | 2798 | $\ldots$ | $<0.20$ |  | $<0.21$ |  | 1.76 | 0.38 | ... |  | ... |
| ESO 163-G011. | 073659.4 | -550438 | I | 2822 | 11.9 | 0.37 | 0.11 | 0.63 | 0.15 | 5.32 | 1.65 | 14.64 | 1.05 | 10.20 |
| ESO 432-IG006 | 084225.6 | -313048 | I | 4846 | 16.4 | 0.88 | 0.09 | 1.14 | 0.11 | 7.15 | 0.73 | 9.91 | 1.00 | 10.67 |
| ESO 564-G010 | 090029.0 | -203100 | O | ... | 14.7 | $<0.20$ |  | $<0.30$ | ... | ... | $\ldots$ | $\ldots$ |  | ... |
| ESO 564-G011. | 090030.0 | -203136 | I | 2596 | 14.5 | 0.38 | 0.05 | 1.45 | 0.17 | 9.17 | 1.00 | 11.85 | 1.20 |  |
| IRAS 0911-1006 e. | 091110.7 | -100703 | O | 16449 | ... | $<0.20$ | $\ldots$ | 0.67 | 0.07 | 7.39 | 0.70 | 13.31 | 1.40 | 11.79 |
| IRAS 0911-1006 w.. | 091113.0 | -10 0654 | O | 16449 |  | ... |  | ... | ... | $\ldots$ | $\ldots$ | ... |  |  |
| NGC $2798=$ Arp 283 | 091410.0 | 421235 | I | 1739 | 13.0 | 0.87 | 0.09 | 3.20 | 0.20 | 19.29 | 0.20 | 27.98 | 0.96 | 10.22 |
| NGC 2799. | 091417.7 | 421214 | O | 1755 | 14.3 | $<0.30$ |  | $<0.40$ | ... | <1.50 |  | $<1.00$ |  |  |
| NGC 2854. | 092038.5 | 492515 | 1 | 2741 | 13.8 | 0.20 | 0.05 | 0.41 | 0.10 | 2.49 | 0.25 | 5.29 | 0.42 | 9.79 |
| NGC 2856................................................ | 092054.6 | 492748 | 1 | 2638 | 14.1 | 0.43 | 0.04 | 1.24 | 0.07 | 6.27 | 0.69 | 11.57 | 0.93 | 10.13 |
| NGC 2857. | 092113.3 | 493436 | I | 4887 | 12.9 | 0.15 | 0.03 | 0.22 | 0.08 | 0.91 | 0.10 | 2.78 | 0.25 | 9.94 |
| MCG +08-18-012 | 093312.0 | 484141 | O | ... | 15.0 | $<0.10$ | ... | $<0.24$ | ... | <1.0 | ... | ... | ... | . . |
| MCG +08-18-013 . | 093318.8 | 484157 | I | 7777 | 15.0 | $<0.20$ | $\ldots$ | 0.79 | 0.24 | 6.23 | 0.45 | 9.34 | 0.67 | 11.03 |
| NGC $2992=$ Arp 245 | 094317.3 | -140839 | I | 2314 | 13.1 | 0.74 | 0.07 | 1.57 | 0.17 | 7.34 | 0.80 | $11.60{ }^{\text {b }}$ | 3.00 | 10.06 |
| NGC 2993. | 094323.6 | -140812 | I | 2420 | 13.1 | 0.58 | 0.06 | 1.88 | 0.20 | 10.92 | 1.10 | $14.64{ }^{\text {b }}$ | 3.80 | 10.25 |
| IC 563 = Arp 303. | 094344.6 | 031634 | I | 6093 | 14.8 | 0.38 | 0.05 | $<0.18$ | ... | $2.41{ }^{\text {b }}$ | 0.20 | $5.87{ }^{\text {b }}$ | 0.60 | 10.50 |
| IC 564 | 094345.3 | 031807 | I | 6056 | 14.1 | 0.30 | 0.05 | 0.24 | 0.04 | $3.07{ }^{\text {b }}$ | 0.25 | $6.56{ }^{\text {b }}$ | 0.70 | 10.57 |
| IC 2522 | 095257.1 | -32 5355 | I | 3012 | 12.6 | 0.89 | 0.23 | 0.94 | 0.13 | 5.15 | 0.22 | 13.32 | 1.24 | 10.23 |
| IC 2523 | 095258.9 | -32 5817 | I | 2611 | 13.6 | $<0.54$ | 0.23 | 0.32 | 0.13 | 2.74 | 0.22 | 6.03 | 1.01 | 9.80 |
| IC $2545{ }^{\text {c }}$. | 100352.2 | -33 3823 | I | 10230 | 15.0 | 0.38 | 0.11 | 1.29 | 0.22 | 9.71 | 0.34 | 9.60 | 0.46 | 11.41 |
| NGC 3165. | 101055.8 | 033725 | O | 1332 | 14.5 | $<0.15$ |  | $<0.30$ | ... | $<0.10$ |  | $<0.55$ | ... | $\ldots$ |
| NGC 3166. | 101109.3 | 034025 | I | 1345 | 11.3 | 0.51 | 0.30 | 0.76 | 0.06 | 5.98 | 0.04 | 12.57 | 0.14 | 9.55 |
| NGC 3169. | 101138.7 | 034303 | I | 1233 | 11.1 | 1.31 | 0.04 | 1.74 | 0.04 | 8.24 | 0.03 | 23.59 | 0.12 | 9.68 |
| NGC $3226=$ UGC 5617. | 102043.2 | 200906 | O | 1151 | 12.3 | $<0.27$ | . | $<0.18$ | $\ldots$ | <0.60 | .. | $\ldots$ | ... | $\ldots$ |
| NGC 3227 = UGC 5620........................... | 102047.3 | 200703 | I | 1157 | 11.1 | 1.11 | 0.12 | 2.04 | 0.21 | 9.01 | 1.00 | 19.11 | 2.00 | 9.60 |
| NGC 3262. | 102657.5 | -43 5413 | O | 2864 | 14.2 | $<0.35$ |  | <0.42 | ... | 0.25 |  | $<0.6$ | ... | ... |
| NGC 3263. | 102704.8 | -435154 | 1 | 2842 | 12.5 | 0.86 | 0.19 | 0.93 | 0.22 | 8.88 | 0.36 | 17.30 | 1.03 | 10.36 |
| NGC 3395 = UGC $5931=$ Arp 270........... | 104702.6 | 331444 | I | 1620 | 12.4 | 0.28 | 0.05 | 0.72 | 0.08 | $6.79{ }^{\text {b }}$ | 2.10 | $13.00^{\text {b }}$ | 3.00 | 9.75 |
| NGC 3396................................................ | 104708.9 | 331518 | I | 1625 | 12.6 | 0.23 | 0.04 | 0.98 | 0.11 | $6.15{ }^{\text {b }}$ | 2.10 | $9.56{ }^{\text {b }}$ | 3.00 | 9.67 |

TABLE 1-Continued

| Name <br> (1) | $\begin{gathered} \text { R.A. } \\ \text { (B1950) } \end{gathered}$ (2) | Decl. (B1950) (3) | Position Type <br> (4) | $\begin{gathered} c z \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \\ (5) \end{gathered}$ | $\begin{aligned} & m_{z} \\ & (6) \end{aligned}$ | $12 \mu \mathrm{~m}$ (Jy) (7) | $\sigma_{12}$ <br> (Jy) <br> (8) | $25 \mu \mathrm{~m}$ (Jy) (9) | $\sigma_{25}$ <br> (Jy) <br> (10) | $60 \mu \mathrm{~m}$ (Jy) <br> (11) | $\sigma_{60}$ $(\mathrm{Jy})$ <br> (12) | $100 \mu \mathrm{~m}$ (Jy) (13) | $\sigma_{100}$ <br> (Jy) <br> (14) | $L_{\text {FIR }}$ <br> $\left(L_{\odot}\right)$ <br> (15) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NGC 3413. | 104834.0 | 330156 | I | 645 | 13.1 | $<0.10$ | .. | $<0.17$ | ... | 1.37 | 0.40 | 2.15 | 0.10 | 8.22 |
| NGC 3424. | 104900.1 | 330954 | I | 1501 | 13.2 | 0.59 | 0.04 | 0.94 | 0.04 | 9.03 | 0.09 | 17.03 | 0.21 | 9.81 |
| NGC 3430. | 104925.2 | 331303 | I | 1585 | 12.2 | 0.38 | 0.08 | 0.78 | 0.05 | 4.36 | 0.07 | 10.88 | 1.00 | 9.59 |
| UGC 6016.. | 105112.1 | 543313 | O | 1493 | 17.0 | $<0.10$ | ... | $<0.15$ | ... | $<0.50$ | ... | $<0.50$ | $\ldots$ | ... |
| NGC 3448. | 105139.6 | 543427 | I | 1350 | 12.5 | 0.34 | 0.11 | 0.76 | 0.21 | 6.74 | 0.35 | 12.17 | 0.47 | 9.58 |
| ESO 264-G057 | 105645.8 | -431026 | I | 5156 | 15.0 | 0.37 | 0.10 | 0.95 | 0.23 | 6.78 | 0.34 | 17.10 | 3.20 | 10.81 |
| ESO 264-G058 | 105648.2 | -430600 | I | 6120 | 17.0 | $<0.30$ | ... | $<0.30$ | ... | 0.61 | 0.18 | $<1.50$ | ... | ... |
| NGC 3511. | 110057.0 | -22 4900 | I | 1106 | 11.5 | 0.92 | 0.11 | 1.16 | 0.13 | 9.34 | 0.13 | 23.93 | 0.41 | 9.62 |
| NGC 3513. | 110119.2 | -22 5828 | I | 1194 | 11.9 | $<0.12$ | ... | $<0.48$ |  | 3.27 | 0.06 | 7.69 | 0.23 | 9.21 |
| MCG +07-23-019 = VV32. | 110106.8 | 410712 | I | 10356 | ... | $<0.30$ | $\ldots$ | 0.42 | 0.13 | 6.57 | 0.19 | 11.38 | 0.44 | 11.33 |
| IC $2810=$ UGC 6436a. | 112308.4 | 145708 | O | 10243 | 14.9 | $<0.25$ | $\ldots$ | $<0.35$ | ... | $5.10{ }^{\text {b }}$ | 1.00 | 10.98 | 0.30 | ... |
| IC $2810 \mathrm{~b}=$ UGC 6436b. | 112313.0 | 145639 | O | 10240 | 15.4 | $<0.25$ |  | $<0.41$ |  | $2.10{ }^{\text {b }}$ | 1.00 | ... | ... |  |
| NGC $3690=$ Arp 299 | 112542.0 | 585017 | I | 3131 | 11.8 | 3.90 | 0.40 | 24.14 | 2.40 | 121.64 | 12.50 | 122.45 | 12.50 | 11.48 |
| NGC 3893... | 114601.8 | 485913 | I | 973 | 11.2 | 1.7 | 0.19 | 2.09 | 0.33 | 15.76 | 0.17 | 38.38 | 0.53 | 9.72 |
| NGC 3896. | 114626.8 | 490139 | O | 869 | 13.9 | $<0.21$ | ... | $<0.18$ |  | $<0.33$ |  | $<0.30$ |  |  |
| NGC $3991=$ Arp 313 | 115454.0 | 323600 | I | 3111 | 14.2 | $<0.12$ | . $\cdot$. | 0.21 | 0.04 | 2.92 | 0.28 | 4.22 | 0.45 | 9.90 |
| NGC 3994.. | 115502.3 | 323323 | I | 3096 | 13.3 | 0.32 | 0.04 | 0.46 | 0.05 | 4.98 | 0.50 | 10.31 | 1.20 | 10.19 |
| NGC 3995. | 115510.3 | 323424 | I | 3254 | 12.7 | $<0.18$ | ... | 0.64 | 0.07 | 3.75 | 0.30 | 6.63 | 0.68 | 10.08 |
| NGC 4038 ${ }^{\text {a }}$ | 115920.0 | -183549 | I | 1642 | 10.9 | 2.92 | 0.38 | 7.11 | 1.13 | 46.88 | 1.70 | 85.69 | 8.20 | 10.59 |
| ESO 440-IG058 = VV 835. | 120417.3 | -314017 | I | 6818 | 16.0 | 0.37 | 0.06 | 0.97 | 0.09 | 7.30 | 0.78 | 13.40 | 1.50 | 11.02 |
| NGC 4169.. | 120947.0 | 292730 | O | 3784 | 13.2 | $<0.35$ | ... | $<0.30$ |  | $<0.33$ | ... | $<0.80$ | ... | ... |
| NGC 4170.. | 120950.2 | 292857 | O | 1127 | 13.6 | $<0.30$ | $\ldots$ | $<0.10$ | $\ldots$ | $<0.30$ | $\ldots$ | $<0.80$ | $\ldots$ | $\ldots$ |
| NGC 4174.. | 120954.9 | 292536 | O | 3980 | 14.3 | $<0.24$ | $\ldots$ | $<0.33$ |  | $<0.21$ | ... | $<1.20$ | ... |  |
| NGC 4175. | 120958.6 | 292652 | I | 3956 | 14.2 | 0.40 | 0.13 | 0.77 | 0.26 | 5.70 | 0.31 | 10.58 | 0.61 | 10.44 |
| ESO 267-G029 | 121115.0 | -46 5947 | I | 5445 | 14.2 | 0.23 | 0.09 | 0.82 | 0.15 | 5.17 | 0.20 | 9.05 | 0.96 | 10.67 |
| ESO 267-G030 | 121134.8 | -465656 | I | 5543 | 14.2 | 0.32 | 0.09 | 0.75 | 0.13 | 4.73 | 0.18 | 10.91 | 0.91 | 10.70 |
| IC 3153. | 121703.8 | 054033 | I | 11646 | 14.8 | $<0.20$ |  | $<0.20$ | $\ldots$ | 0.60 | 0.18 | 1.32 | 0.41 | 10.38 |
| NGC 4266. | 121708.2 | 054906 | I | 2495 | 14.6 | $<0.20$ | $\ldots$ | $<0.30$ | $\ldots$ | 0.31 | 0.1 | 2.09 | 1.31 | 9.10 |
| NGC 4268. | 121713.9 | 053340 | O | 2374 | 13.8 | $<0.20$ | $\ldots$ | $<0.20$ | $\ldots$ | $<0.30$ | ... | $<0.51$ | ... |  |
| NGC 4270. | 121716.2 | 054427 | O | 2357 | 13.1 | $<0.20$ |  | $<0.30$ |  | $<0.50$ |  | $<0.33$ | $\ldots$ |  |
| NGC 4273. | 121723.6 | 053713 | I | 2378 | 12.4 | 1.05 | 0.34 | 1.35 | 0.39 | 11.27 | 0.25 | 21.87 | 1.31 | 10.31 |
| NGC 4277.. | 121730.5 | 053707 | O | 2516 | 13.4 | $<0.20$ | ... | $<0.20$ | ... | $<0.60$ |  | $<1.20$ | $\ldots$ | ... |
| NGC 4281.. | 121748.6 | 053954 | I | 2711 | 12.3 | $<0.20$ | $\ldots$ | $<0.40$ |  | 0.83 | 0.15 | 1.78 | 0.51 | 9.31 |
| NGC $4485=$ Arp 269 | 122803.3 | 415826 | I | 493 | 12.3 | $<0.90$ | . $\cdot$ | 0.53 | 0.10 | 2.16 | 0.70 | 3.83 | 1.05 | 8.21 |
| NGC 4490.. | 122808.9 | 415526 | I | 578 | 10.2 | 2.74 | 0.14 | 5.34 | 0.80 | 50.86 | 6.00 | 88.29 | 1.05 | 9.71 |
| NGC 4567/8 ${ }^{\text {a }}$ | 123402.3 | 113106 | I | 2260 | 11.1 | 2.1 | 0.25 | 2.93 | 0.74 | 21.11 | 0.25 | 56.77 | 1.25 | 10.60 |
| IC 3639 .. | 123810.2 | -362845 | I | 3285 | 13.0 | 0.82 | 0.10 | 2.87 | 0.30 | 8.27 | 1.00 | 14.90 | 1.20 | ... |
| ESO 381-G009 | 123816.5 | -362705 | I | 3050 | 13.9 | $<0.16$ | ... | 0.32 | 0.10 | 1.17 | 0.15 | ... | ... | $\ldots$ |
| NGC 4627.. | 123933.5 | 325051 | O | 765 | 13.1 | $<0.27$ | ... | $<0.24$ | ... | $<0.50$ | .. | $<1.80$ | $\ldots$ |  |
| NGC 4631. | 123940.7 | 324854 | I | 606 | 09.8 | 6.81 | 0.25 | 11.24 | 0.68 | 99.69 | 3.50 | 193.26 | 4.10 | 10.07 |
| NGC 4647.. | 124101.0 | 115120 | I | 1414 | 11.9 | 0.76 | 0.08 | 1.06 | 0.16 | 6.04 | 0.24 | 17.56 | 0.81 | 9.67 |
| NGC 4649.. | 124108.4 | 114935 | O | 1114 | 9.8 | 0.33 | 0.05 | $<0.32$ | ... | $<0.24$ | ... | <1.05 | ... | ... |
| NGC $4922 \mathrm{a} / \mathrm{b}$.. | 125901.7 | 293500 | I | 7071 | 14.2 | 0.22 | 0.08 | 1.67 | 0.27 | 6.01 | 1.22 | 7.78 | 0.67 | 10.91 |
| MCG -02-33-098. | 125941.6 | -15 2954 | I | 4773 | 14.5 | $<0.33$ | $\ldots$ | 1.95 | 0.39 | 7.82 | 0.23 | 10.29 | 0.85 | 10.69 |
| MCG -02-33-099. | 125949.1 | -153036 | O | 5019 | 18.0 | $<0.10$ | $\ldots$ | $<0.30$ | ... | $<0.20$ | $\ldots$ | $<0.50$ | ... | ... |
| UGC 8335 NED01 = VV250b = Arp238 ... | 131336.3 | 622332 | O | 9453 | 15.0 | $<0.30$ | $\ldots$ | $<0.50$ | ... | ... | ... | ... | $\ldots$ |  |
| UGC 8335 NED02 = VV250a.................... | 131341.4 | 622321 | I | 9313 | 15.0 | 0.49 | 0.05 | 2.08 | 0.20 | 11.45 | 1.30 | 12.90 | 1.30 | 11.41 |

TABLE 1—Continued

| Name <br> (1) | R.A. (B1950) (2) | Decl. (B1950) <br> (3) | Position Type <br> (4) | $\begin{gathered} c z \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \\ (5) \end{gathered}$ | $m_{z}$ <br> (6) | $12 \mu \mathrm{~m}$ (Jy) (7) | $\sigma_{12}$ <br> (Jy) <br> (8) | $25 \mu \mathrm{~m}$ (Jy) (9) | $\sigma_{25}$ (Jy) (10) | $60 \mu \mathrm{~m}$ (Jy) (11) | $\sigma_{60}$ <br> (Jy) <br> (12) | $100 \mu \mathrm{~m}$ | $\sigma_{100}$ <br> (Jy) <br> (14) | $L_{\text {FIR }}$ $\left(L_{\odot}\right)$ <br> (15) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC 879 | 131655.2 | $-271012$ | I | 1969 | 14.0 | 0.23 | 0.05 | 0.52 | 0.08 | 0.74 | 0.08 | $<0.29$ |  |  |
| NGC 5078. | 131704.6 | -27 0851 | I | 2168 | 11.8 | 1.20 | 0.13 | 1.36 | 0.15 | 10.37 | 1.10 | 35.84 | 5.70 | 10.32 |
| MCG -03-34-063 | 131938.3 | $-162650$ | O | 6394 | 15.0 | $<0.45$ |  | $<0.4$ |  | $3.03{ }^{\text {b }}$ | 1.05 | $3.90{ }^{\text {b }}$ | 1.10 | 10.34 |
| MCG -03-34-064 | 131943.6 | -1628 05 | O | 4959 | 14.6 | 0.84 | 0.24 | 3.08 | 0.36 | $4.63{ }^{\text {b }}$ | 1.13 | $2.80{ }^{\text {b }}$ | 1.10 | 10.44 |
| NGC $5216=$ UGC 8528 | 133023.0 | 630132 | I | 2949 | 13.6 | <0.09 |  | $<0.18$ |  | 0.19 | 0.03 | $<0.05$ |  |  |
| NGC $5218=$ UGC 8529 | 133027.1 | 630132 | I | 2860 | 13.1 | 0.36 | 0.03 | 0.92 | 0.08 | 7.14 | 0.70 | 14.38 | 1.44 | 10.28 |
| NGC 5257 = Arp 240 | 133719.7 | 010533 | O | 6798 | 12.9 | $0.52^{\text {b }}$ | 0.16 | $1.18{ }^{\text {b }}$ | 0.30 | $8.10{ }^{\text {b }}$ | 2.00 | $13.63{ }^{\text {b }}$ | 3.00 | 11.05 |
| NGC 5258... | 133724.6 | 010506 | O | 6757 | 12.9 | $0.25{ }^{\text {b }}$ | 0.10 | $0.78{ }^{\text {b }}$ | 0.30 | $3.94{ }^{\text {b }}$ | 1.00 | $7.27^{\text {b }}$ | 1.80 | 10.75 |
| NGC 5331 ${ }^{\text {a }}$ | 134943.4 | 022060 | I | 9906 | 13.7 | $<0.30$ | ... | 0.79 | 0.15 | 6.27 | 0.33 | 10.71 | 0.50 | 11.27 |
| NGC $5394=$ UGC $8898=$ Arp 84.... | 135622.5 | 374200 | I | 3427 | 13.7 | 0.52 | 0.05 | 1.19 | 0.11 | $5.62{ }^{\text {b }}$ | 1.41 | $10.43{ }^{\text {b }}$ | 3.10 | 10.31 |
| NGC 5395. | 135629.3 | 374003 | I | 3487 | 12.1 | 0.40 | 0.04 | 0.48 | 0.06 | $6.86{ }^{\text {b }}$ | 1.50 | $14.21{ }^{\text {b }}$ | 3.10 | 10.44 |
| IC 4356 | 135636.6 | 374359 | O | ... | 16.2 | $<0.20$ | $\ldots$ | $<0.27$ | $\ldots$ | $<0.45$ |  | $<0.10$ |  | ... |
| NGC $5426=$ Arp 271 | 140047.7 | -054947 | I | 2621 | 12.7 | $<0.70$ | $\ldots$ | $<0.90$ |  | 3.30 | 0.37 | 8.58 | 1.52 | 9.92 |
| NGC 5427. | 140048.6 | -054727 | I | 2618 | 11.9 | 0.74 | 0.09 | 0.96 | 0.10 | 7.50 | 0.80 | 16.47 | 3.32 | 10.24 |
| NGC 5506. | 141039.1 | -02 5826 | I | 1815 | 13.4 | 1.48 | 0.09 | 4.29 | 0.15 | 8.82 | 0.10 | 8.96 | 0.12 | 9.87 |
| NGC 5507. | 141043.9 | -02 5455 | O | 1852 | 13.6 | $<0.16$ |  | $<0.28$ |  | $<0.38$ |  | $<0.66$ | ... | ... |
| NGC 5595. | 142128.4 | -162955 | I | 2691 | 13.1 | 0.67 | 0.04 | 0.98 | 0.05 | 8.84 | 0.04 | 17.23 | 0.14 | 10.31 |
| NGC 5597. | 142142.2 | -1632 19 | I | 2619 | 12.6 | 0.63 | 0.04 | 1.85 | 0.06 | 8.90 | 0.06 | 15.30 | 0.16 | 10.27 |
| NGC 5734. | 144219.0 | -20 3936 | I | 4074 | 13.7 | 0.78 | 0.15 | 0.98 | 0.07 | 8.09 | 0.15 | 17.48 | 2.45 | 10.65 |
| NGC 5743. | 144220.0 | -20 4212 | I | 4216 | 13.8 | 0.59 | 0.13 | 0.75 | 0.08 | 5.21 | 0.15 | 9.78 | 1.27 | 10.46 |
| IC 4518 A/B | 145426.0 | -42 5554 | I | 4921 | 15.0 | 0.53 | 0.05 | 1.52 | 0.15 | 9.59 | 1.00 | 22.22 | 2.30 | 10.90 |
| UGC 9618 NED $2=$ VV 340 | 145447.8 | 244905 | I | 10173 | 14.6 | $0.53{ }^{\text {a }}$ | 0.13 | 0.66 | 0.18 | 7.08 | 2.00 | 15.48 | 5.00 | 11.39 |
| UGC 9618 NED 1. | 145448.1 | 244821 | O | 9776 | 15.3 | $\ldots$ |  | $<0.21$ | ... | $<1.50$ |  | <5.00 |  |  |
| NGC 5793. | 145637.1 | -1629 40 | I | 3491 | 14.2 | 0.25 | 0.04 | 0.68 | 0.07 | 6.14 | 0.22 | 10.57 | 0.19 | 10.35 |
| NGC 5796. | 145636.5 | $-162530$ | O | 2962 | 12.7 | $<0.25$ | ... | $<0.30$ | ... | $<0.36$ | ... | $<0.35$ | ... | ... |
| I Zw 107. | 151619.7 | 425536 | I | 12043 | 15.0 | 0.31 | 0.09 | 1.47 | 0.14 | 9.57 | 0.19 | 10.75 | 0.35 | 11.56 |
| NGC 5915. | 151847.7 | -125456 | I | 2291 | 13.0 | 0.74 | 0.04 | 1.84 | 0.05 | 11.71 | 0.05 | 17.45 | 0.22 | 10.25 |
| NGC 5916a. | 151828.9 | -1255 29 | I | 2338 | 14.6 | $<0.60$ | $\ldots$ | $<0.90$ | $\ldots$ | 0.79 | 0.04 | $<1.50$ | ... | ... |
| NGC 5916. | 151852.0 | -125937 | I | 2338 | 14.2 | $<0.15$ | $\ldots$ | $<0.24$ | $\ldots$ | 1.11 | 0.06 | 3.04 | 0.35 | 9.42 |
| NGC 5929. | 152418.9 | 415041 | O | 2561 | 14.1 | $<0.24$ | ... | $<0.30$ | $\ldots$ | $<1.00$ | $\ldots$ | $\ldots$ | ... | $\ldots$ |
| NGC 5930. | 152420.7 | 415100 | O | 2672 | 13.6 | 0.45 | 0.09 | 1.84 | 0.19 | 9.71 | 0.19 | 14.71 | 0.66 | 10.30 |
| NGC 5953. | 153213.7 | 152143 | I | 1965 | 13.3 | 0.72 | 0.12 | 1.85 | 0.11 | 12.00 | 0.16 | 21.18 | 0.85 | 10.15 |
| NGC 5954. | 153215.7 | 152210 | O | 1959 | 13.7 | $<0.15$ |  | $<0.40$ |  |  |  |  |  |  |
| NGC 6285. | 165736.1 | 590153 | I | 5691 | 15.3 | 0.20 | 0.05 | 0.27 | 0.05 | $2.60{ }^{\text {b }}$ | 0.60 | $6.60{ }^{\text {b }}$ | 2.30 | 10.48 |
| NGC 6286. | 165745.0 | 590041 | I | 5501 | 14.1 | 0.45 | 0.06 | 0.55 | 0.06 | $8.03{ }^{\text {b }}$ | 2.00 | $18.30{ }^{\text {b }}$ | 4.00 | 10.92 |
| IC 4686/7 ${ }^{\text {c }}$. | 180919.4 | -574428 | I | 5200 | 13.8 | 0.94 | 0.20 | 2.96 | 0.20 | $16.08^{\text {b }}$ | 1.35 | 28.55 | 1.09 | 11.12 |
| IC 4689 . | 180922.1 | -574545 | I | 4949 | 15.0 | $<0.25$ | ... | 0.71 | 0.20 | $5.00^{\text {b }}$ | 1.20 | ... | ... | ... |
| NGC $6621=\operatorname{Arp} 81$. | 181309.3 | 682050 | I | 6284 | 14.0 | 0.31 | 0.05 | 1.02 | 0.10 | 7.02 | 1.20 | 12.19 | 2.30 | 10.93 |
| NGC 6622. | 181314.4 | 682015 | O | 6230 | 16.0 | $<0.12$ | ... | $<0.30$ | . | $<1.00$ |  | $<1.30$ | ... | ... |
| CGCG 142-033 = Zw 142. | 181424.0 | 220500 | O | 5353 | 15.6 | $<0.20$ | $\ldots$ | $<0.30$ |  | $<0.30$ |  | $<1.00$ |  | $\ldots$ |
| CGCG 142-034. | 181432.6 | 220536 | I | 5599 | 15.6 | 0.27 | 0.06 | 0.64 | 0.08 | 6.55 | 0.80 | 13.25 | 1.60 | 10.82 |
| NGC 6670. | 183256.4 | 595058 | I | 8650 | ... | 0.47 | 0.05 | 1.10 | 0.10 | 9.24 | 0.90 | 14.84 | 1.50 | 11.31 |
| CGCG 301-032. | 183332.8 | 594914 | I | 8699 | 15.2 | 0.15 | 0.04 | 0.15 | 0.04 | 0.90 | 0.10 | 1.64 | 0.18 | 10.32 |
| NGC 6745. | 190003.3 | 404023 | I | 4545 | 13.3 | 0.45 | 0.05 | 0.92 | 0.12 | 6.84 | 0.70 | 13.01 | 1.40 | 10.65 |
| NGC $6786=$ VV 414. | 191152.9 | 731931 | I | 7510 | 13.8 | 0.15 | 0.04 | 0.55 | 0.05 | $3.40^{\text {b }}$ | 0.60 | $6.11{ }^{\text {b }}$ | 1.40 | 10.77 |
| UGC 11415............................................. | 191204.0 | 732027 | I | 7555 | 14.9 | 0.21 | 0.04 | 0.91 | 0.09 | $4.51{ }^{\text {b }}$ | 0.84 | $5.89{ }^{\text {b }}$ | 1.40 | 10.85 |
| NGC 6921............................................... | 202620.8 | 253323 | O | 4391 | 14.4 | $<0.25$ | ... | 0.62 | 0.20 | 3.97 | 1.20 | 7.48 | 2.30 | 10.38 |
| MCG +04-48-002 ..................................... | 202626.8 | 253400 | O | 4259 | 18.0 | 0.51 | 0.09 | 0.71 | 0.20 | 8.15 | 1.20 | 12.50 | 2.70 | 10.63 |

TABLE 1—Continued

| Name <br> (1) | $\begin{aligned} & \text { R.A. } \\ & \text { (B1950) } \\ & (2) \end{aligned}$ | Decl. (B1950) (3) | Position Type <br> (4) | $\begin{gathered} c z \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \\ (5) \end{gathered}$ | $m_{z}$ <br> (6) | $12 \mu \mathrm{~m}$ (Jy) (7) | $\sigma_{12}$ <br> (Jy) <br> (8) | $\begin{gathered} 25 \mu \mathrm{~m} \\ \text { (Jy) } \\ \text { (9) } \end{gathered}$ | $\begin{gathered} \sigma_{25} \\ (\mathrm{Jy}) \\ (10) \end{gathered}$ | $60 \mu \mathrm{~m}$ (Jy) (11) | $\sigma_{60}$ <br> (Jy) <br> (12) | $100 \mu \mathrm{~m}$ (Jy) (13) | $\sigma_{100}$ <br> (Jy) <br> (14) | $L_{\text {FIR }}$ $\left(L_{\odot}\right)$ (15) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NGC 6926. | 203029.8 | -02 1201 | I | 5970 | 13.2 | 0.66 | 0.11 | 0.91 | 0.18 | 7.85 | 0.33 | 18.50 | 1.64 | 10.99 |
| NGC 6929. | 203045.9 | -02 1231 | O | 6174 | 14.4 | $<0.50$ | ... | $<0.40$ |  | $<0.50$ |  | $<0.50$ |  |  |
| ESO 286-IG033 | 210050.6 | -43 4359 | I | 4990 | ... | . ${ }^{\text {d }}$ | $\ldots$ | 0.43 | 0.08 | 1.33 | 0.10 | 1.42 | 0.17 | 9.93 |
| ESO 286-IG035 | 210052.7 | -43 4731 | I | 5208 | 14.7 | d | ... | 1.12 | 0.14 | 8.17 | 0.90 | 14.16 | 1.60 | 10.83 |
| NGC 7172. | 215907.1 | -320642 | I | 2575 | 12.9 | 0.72 | 0.06 | 1.06 | 0.11 | 6.06 | 0.61 | 15.10 | 1.35 | 10.16 |
| NGC 7174/6. | 215911.9 | -32 1403 | I | 2778 | 12.2 | 0.25 | 0.04 | 0.61 | 0.06 | 3.65 | 0.40 | 10.15 | 1.00 | 10.03 |
| NGC $7253=\operatorname{Arp} 278{ }^{\text {a }}$ | 221711.7 | 290841 | I | 4718 | 14.4 | 0.43 | 0.05 | 0.70 | 0.09 | 6.38 | 0.70 | 13.63 | 1.50 | 10.67 |
| NGC 7469. | 230044.4 | 083619 | I | 4916 | 13.0 | 1.63 | 0.35 | 5.70 | 0.19 | 23.13 | 0.64 | 39.91 | 1.36 | 11.22 |
| IC 5283 | 230047.0 | 083726 | O | 4894 | 14.8 | $<0.30$ | ... | $<0.31$ | ... | ... |  | ... |  | ... |
| NGC 7537. | 231200.9 | 041321 | I | 2674 | 13.9 | $<0.45$ | $\ldots$ | $<0.60$ | $\ldots$ | 1.25 | 0.18 | 1.67 | 0.64 | 9.39 |
| NGC 7541. | 231211.8 | 041536 | I | 2678 | 12.4 | 1.34 | 0.24 | 2.36 | 0.32 | 20.81 | 0.75 | 43.95 | 1.59 | 10.69 |
| NGC 7552. | 231325.4 | -425127 | I | 1585 | 11.3 | 3.49 | 0.36 | 12.43 | 1.40 | 78.38 | 8.00 | 104.85 | 12.00 | 10.73 |
| NGC 7582. | 231538.3 | -423154 | I | 1575 | 11.4 | 2.54 | 0.30 | 9.43 | 1.10 | 52.25 | 5.60 | 87.60 | 9.30 | 10.59 |
| NGC 7590.. | 231610.3 | -42 3045 | I | 1596 | 12.1 | 1.02 | 0.10 | 1.23 | 0.14 | 8.54 | 0.90 | 19.57 | 2.40 | 9.87 |
| NGC 7599.. | 231635.9 | -42 3154 | I | 1654 | 12.1 | 1.06 | 0.10 | 1.40 | 0.18 | 6.95 | 1.40 | 19.45 | 2.00 | 9.86 |
| NGC 7592. | 231547.4 | -04 4126 | I | 7328 | 11.4 | 0.62 | 0.19 | 1.26 | 0.30 | 7.81 | 0.36 | 10.66 | 0.99 | 11.07 |
| ESO 077-IG 014. | 231758.5 | -69 2928 | I | 11400 | 16.5 | $<0.23$ | ... | 0.61 | 0.17 | 6.02 | 0.35 | 10.62 | 1.41 | 11.38 |
| NGC 7714 = UGC 12699 | 233341.0 | 015234 | I | 2798 | 13.0 | 0.56 | 0.05 | 3.15 | 0.30 | 10.73 | 1.10 | 12.46 | 1.30 | 10.34 |
| NGC 7715.. | 233348.5 | 015248 | O | 2770 | 14.7 | $<0.15$ | $\ldots$ | $<0.21$ | ... | $<0.15$ |  | $<0.40$ | ... | ... |
| MCG -01-60-021 | 233911.6 | -03 5646 | I | 6595 | 16.5 | $<0.27$ | $\ldots$ | $<0.20$ | ... | 0.81 | 0.31 | 1.42 | 0.31 | 10.03 |
| MCG -01-60-022 | 233927.1 | -03 5333 | I | 6966 | 14.5 | 0.37 | 0.13 | 0.78 | 0.08 | 5.11 | 0.32 | 9.45 | 0.31 | 10.89 |
| NGC 7752. | 234427.1 | 291052 | O | 5072 | 15.0 | 0.64 | 0.05 | 0.54 | 0.15 | 3.96 | 1.31 | 8.02 | 2.41 | 10.52 |
| NGC 7753. | 234433.3 | 291221 | O | 5163 | 12.8 | ... | ... | 0.34 | 0.15 | 2.68 | 0.88 | 8.87 | 2.30 | 10.47 |
| NGC 7769.. | 234832.2 | 195223 | I | 4214 | 12.8 | 0.52 | 0.04 | 0.97 | 0.08 | 5.21 | 0.40 | 13.58 | 1.40 | 10.53 |
| NGC 7771.. | 234851.7 | 194952 | I | 4287 | 13.1 | 1.23 | 0.16 | 2.90 | 0.22 | 20.93 | 1.78 | 44.85 | 3.80 | 11.11 |
| UGC 12812. | 234845.8 | 201800 | O | 5326 | 15.5 | $<0.25$ | ... | $<0.30$ | ... | <1.00 |  | $<1.30$ | ... | ... |
| Mrk 331 = MCG +03-60-036 .................... | 234853.6 | 201824 | I | 5541 | 14.9 | 0.87 | 0.13 | 3.02 | 0.24 | 18.43 | 1.60 | 22.56 | 2.50 | 11.18 |
| UGC 12914............................................. | 235904.0 | 231223 | O | 4371 | 13.1 | $<0.14$ | ... | $<0.60$ | ... | 2.52 | 0.70 | 4.05 | 1.90 | 10.15 |
| UGC 12915............................................. | 235908.6 | 231259 | O | 4336 | 14.0 | 0.36 | 0.06 | 1.16 | 0.24 | 5.18 | 1.70 | 13.05 | 2.80 | 10.54 |



 edition of the Astronomical Journal.
${ }^{a}$ Flux lies between galaxies.
${ }^{\mathrm{b}}$ Resolved using component fitting.
c Except NGC 7319/7320.
${ }^{\mathrm{d}}$ Bad scan.

TABLE 2
Integrated Flux Densities of RBGS Interacting Galaxies Dropped from the RBGS

| Name (1) | $\begin{aligned} & \text { R.A. } \\ & \text { (B1950) } \\ & \text { (2) } \end{aligned}$ | Decl. (B1950) (3) | Position Type <br> (4) | $\begin{gathered} c z \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \\ (5) \end{gathered}$ | $\begin{aligned} & m_{z} \\ & (6) \end{aligned}$ | $\begin{gathered} 12 \mu \mathrm{~m} \\ (\mathrm{Jy}) \\ \text { (7) } \end{gathered}$ | $\sigma_{12}$ (Jy) (8) | $\begin{gathered} 25 \mu \mathrm{~m} \\ (\mathrm{Jy}) \\ (9) \end{gathered}$ | $\begin{aligned} & \sigma_{25} \\ & (\mathrm{Jy}) \\ & (10) \end{aligned}$ | $\begin{gathered} 60 \mu \mathrm{~m} \\ (\mathrm{Jy}) \\ (11) \end{gathered}$ | $\begin{aligned} & \sigma_{60} \\ & (\mathrm{Jy}) \\ & (12) \end{aligned}$ | $100 \mu \mathrm{~m}$ (Jy) (13) | $\begin{aligned} & \sigma_{100} \\ & (\mathrm{Jy}) \\ & (14) \end{aligned}$ | $\begin{gathered} L_{\mathrm{FIR}} \\ \left(L_{\odot}\right) \\ (15) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NGC 274. | 004830.0 | -07 1945 | O | 1750 | 12.8 | $<0.40$ | $\ldots$ | $<0.25$ |  | $<1.20$ |  |  |  | $\ldots$ |
| NGC 275. | 004833.0 | -07 2009 | I | 1750 | 13.2 | $<0.40$ |  | 0.91 | 0.08 | 5.38 | 0.40 | 9.54 | 1.15 | 9.70 |
| NGC $1143=$ UGC $2388=\operatorname{Arp} 118 . \ldots . .$. | 025236.2 | -00 2247 | O | 8459 | 13.2 | $<0.30$ | $\ldots$ | <0.10 |  | <1.10 | .. | <1.50 |  |  |
| NGC 1144 = UGC 2389 . | 025238.6 | -00 2311 | I | 8647 | 13.8 | 0.44 | 0.05 | 0.85 | 0.08 | 5.70 | 0.68 | 11.75 | 1.72 | 11.14 |
| UGC 3094.. | 043239.0 | 190412 | I | 7408 | 16.5 | 0.60 | 0.06 | 1.10 | 0.10 | 7.16 | 0.72 | 16.39 | 1.70 | 11.13 |
| 2MASX J04354305+1909568.............. | 043247.8 | 190351 | O |  |  | $<0.12$ | ... | $<0.22$ |  | $<0.28$ |  | $<2.00$ |  |  |
| IRAS 15335-0513.. | 153332.6 | -05 1355 | I | 8186 | 16.7 | $<0.23$ | $\ldots$ | 0.88 | 0.10 | 5.83 | 0.60 | 10.10 | 1.00 | 11.07 |
| 2MASX J15360897-0521513 .............. | 153608.9 | -05 2152 | O | 7077 | 16.6 | $<0.12$ | $\ldots$ | $<0.15$ |  | $<0.25$ | $\ldots$ | $<0.30$ | $\ldots$ | . |
| NGC 6907/8....................................... | 202208.0 | -24 5817 | I | 3161 | 11.9 | 1.29 | 0.13 | 2.42 | 0.25 | 14.76 | 1.50 | 31.78 | 3.50 | 10.69 |

Note.-When a target was unresolved, the integrated HIRES flux (sum of all galaxy components) is given for the component dominant at other wavelengths, with the remaining components having no data. Upper limits are indicated with a < symbol, with the associated $1 \sigma$ values blank. See text for detailed explanation of columns.


Fig. 2.-Illustration of amplification of high-sigma noise outliers by the HIRES process, resulting in noise "spikes" similar in size to the reconstructed IRAS beam and having total fluxes of $\approx 0.1 \mathrm{Jy}$. This is the $60 \mu \mathrm{~m}$ image of Arp 271.

Upper limits are denoted with a " $<$ " and are the flux measured at the known optical location of the galaxy in an aperture that has the same size as the effective IRAS beam.

Column (15).-Log of the far-IR luminosity, in units of solar luminosities. This is the luminosity from $40-122 \mu \mathrm{~m}$ (Helou et al. 1988). We assume $H_{0}=75 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}$. This quantity is useful for photometric study, is not very sensitive to the shape of the spectral energy distribution, and is the quantity tabulated by Surace et al. (1993). The $L_{\text {FIR }}$ described in the RBGS is the luminosity from $1-500 \mu \mathrm{~m}$. It was derived by applying a correction factor to the flux between $40-122 \mu \mathrm{~m}$ based on the $60 / 100 \mu \mathrm{~m}$ color (Lonsdale \& Helou 1985). For the galaxies described here, the median correction factor is $1.44 \pm 0.09$. In other words, $\log L_{1-500 \mu \mathrm{~m}}=\log L_{40-122 \mu \mathrm{~m}}+0.16$. This is also different from the quantity $L_{\mathrm{IR}}$ described by Sanders \& Mirabel (1996), which is the flux from $8-1000 \mu \mathrm{~m}$, but which generally cannot be computed here since it requires detections in all four $\operatorname{IRAS}$ bands.

### 2.3. Photometric Uncertainties

Evaluating the photometric uncertainty of the HIRES data product is quite difficult. In general, uncertainties arise from three sources, all of which vary in importance depending on the particular field. The integrated flux density uncertainties quoted in Tables 1 and 2 contain measurement and confusion errors but not systematic effects in the overall calibration.

Confusion is the first limitation. The dominant source of noise in HIRES is not photometric background noise but confusion due to noise spike amplification. High-sigma noise peaks are amplified by the deconvolution process; they appear similar to weak point sources with a signal strength as high as 0.1 Jy. These spikes are illustrated in Figure 2. This results in a highly non-Gaussian single-sided noise distribution on spatial scales similar to the beam size, not the pixel size. In those cases in which the galaxy fluxes are less than 0.3 Jy , it becomes difficult to differentiate the target from amplified noise. As a result, quoted upper limits are often quite high, as
this upper limit is set by the flux contained in these noise peaks. Our achieved sensitivity is thus around $0.25-0.3 \mathrm{Jy}$, depending on the wavelength and field geometry. Similarly, quoted uncertainties are often also high, depending on the amplitude of these spikes. Apertures similar to the effective beam size were used to evaluate a median false signal due to the noise spikes. These spikes are the dominant source of uncertainty for faint objects.

The technique used to derive the photometry is the second contributor to the photometric uncertainty. In cases in which the galaxies are well separated and aperture photometry could be used, this typically contributes only a few percent to the total error. In those cases in which the galaxies were not well separated and Gaussian fitting was used, this becomes the dominant source of error and can range anywhere from $20 \%-$ $50 \%$ depending on the degree of resolution of the targets.

Absolute photometric calibration is the third major source of uncertainty and is not included in Table 1. As noted in Surace et al. (1993), there are certain caveats to the photometric calibration of the HIRES data product. In particular, the calibration of the IRAS data partly depends on factors such as detector responsivity and dwell time. Known as the $\mathrm{AC} / \mathrm{DC}$ correction, it is the difference in responsivity for point sources versus extremely extended sources, which was characterized as a function of detector dwell time based on the nominal survey slew speed. This is well known for point sources, and hence the Point Source Catalog (PSC) is properly calibrated (Beichmann et al. 1988). However, it is slightly different for small extended sources and is a poorly understood function of source extension. This was seen during the data analysis presented in Paper I, in which it was found that the majority of the HIRES fluxes were significantly greater than the values estimated using one-dimensional coaddition with the ADDSCAN/SCANPI processing available at IPAC. In Paper I this issue was wholly avoided by forcing all of the data onto the same flux scale as the PSC by using the component flux ratios indicated by HIRES to divide up the flux indicated by the BGS, which was produced using the ADDSCAN process, which is known to have the same photometric scale as the PSC.

Figures 3-6 show the difference in flux estimates between this paper and the ADDSCAN/SCANPI values published in the RBGS (Sanders et al. 2003). The data points shown in these figures are limited to cases in which the flux referred to in the RBGS was unambiguous. These are primarily systems that were either unresolved by HIRES (and hence both catalogs have single fluxes) or were sufficiently separated as to have been resolved by both ADDSCAN/SCANPI and HIRES (e.g., NGC 875). The mean offsets between the two catalogs are $27.5,12.8,4.5$, and $5.5 \%$ at $12,25,60$, and $100 \mu \mathrm{~m}$, respectively. The observed scatter around the mean is similar to the estimated flux uncertainties in Table 1. These are particularly significant in the faint $12 \mu \mathrm{~m}$ channel, in which although the minimum requirement for reporting is $\mathrm{S} / \mathrm{N}>3$, the average detection only has $\mathrm{S} / \mathrm{N} \approx 5$. The mean offsets are also similar in size to the observed scatter. Testing of HIRES has shown that measured integrated fluxes of unresolved point sources have an intrinsic scatter of about $8 \%-12 \%$ compared to those of the PSC (Laughlin et al. 1990). In all cases the $1 \sigma$ scatter in offsets observed for the galaxies is larger than the value of the systematic offset. There is no statistically significant trend as a function of flux. In several cases the statistically significant outliers seen in the brighter channels are a result of differences in background estimation between the


FIg. 3.-Fractional difference between one-dimensional ADDSCAN/ SCANPI fluxes (Sanders et al. 2003) and HIRES fluxes at $12 \mu \mathrm{~m}$. The horizontal axis is the RBGS flux in Jy. The vertical axis scaling is such that 0.2 indicates a difference of $20 \%$.
one-dimensional ADDSCAN results and the two-dimensional HIRES results.

While the version of HIRES used in this paper produces data believed to be on the AC scale, appropriate for point sources, it is clear that there are systematic offsets relative to other AC-calibrated IRAS data products. Previous experiments in Paper I showed that the HIRES data product correctly reproduces the photometry of point sources in accordance with the PSC. As the exact source of this offset remains unclear, as does the calibration for small extended sources, this data have not been forced to agree with the RBGS, unlike Surace et al. (1993). This is the source of the variations between the fluxes in Paper I and this work.

As a result of this offset, the results presented here differ from those based solely on PSC-calibrated products by small amounts. The ratio of 60 to $100 \mu \mathrm{~m}$ flux remains unchanged, as the offset is the same in both bands. The infrared luminosities are $5 \%$ higher, a value considerably less than the typical uncertainty. The log of the 12 to $25 \mu \mathrm{~m}$ ratio differs by being 0.05 higher. When appropriate, these offsets will be discussed in $\S 3$.

## 3. RESULTS

### 3.1. Far-IR Properties

The cumulative distribution functions (CDFs) of $L_{\text {FIR }}$, $\log \left(f_{12} / f_{25}\right)$, and $\log \left(f_{60} / f_{100}\right)$ are given in Figures 7-9.


Fig. 4.-Same as Fig. 3, but for $25 \mu \mathrm{~m}$.


Fig. 5.-Same as Fig. 3, but for $60 \mu \mathrm{~m}$.

These distributions only include the galaxies actually detected and resolved by HIRES. These distributions are nearly identical to Surace et al. (1993), which is to be expected, since the shape of the CDF remains unchanged so long as the nature of the incompleteness in the data is random. Since Paper I differed from this paper in being drawn from a parent sample different from the RBGS primarily in spatial extent on the sky, the CDFs are expected to remain the same.

As in Surace et al. (1993), a comparison sample of galaxies was constructed by selecting a subsample drawn from the BGS that had no visible signs of interaction and were not in close pairs (this is the same sample described in Paper I). From these isolated BGS galaxies, we selected a subsample so as to have the same distribution of blue magnitudes as the RBGS close pairs. We can therefore compare the far-IR properties of a far-IR flux-limited sample of interacting pairs to a similarly flux-limited sample of isolated galaxies with the same distribution of optical luminosities.

The CDF for $\log L_{\mathrm{FIR}}$, which is computed from the 60 and $100 \mu$ m fluxes, is shown in Figure 7. The median value of $L_{\text {FIR }}$ is $10^{10.50} L_{\odot}$ for individual, resolved galaxies in the paired and multiple RBGS systems studied here. This is somewhat higher than found in Paper I ( $L_{\text {FIR }}=10^{10.30} L_{\odot}$ ) and cannot be readily attributed to the offsets in calibration, which are of order $5 \%$ at these wavelengths. Kolmogorov-Smirnov statistics indicate that the isolated and paired samples are not drawn from the same parent sample at better than the $99.99 \%$ confidence level. In separated galaxy pairs, then, the interaction process enhances $L_{\text {FIR }}$ by a factor of roughly 3 .


Fig. 6.-Same as Fig. 3, but for $100 \mu \mathrm{~m}$.


FIG. 7.- $L_{\text {FIR }}$ cumulative distribution function for single, isolated galaxies in the RBGS and individual galaxies in multiple systems resolved by HIRES. The resolved individual galaxies in RBGS paired systems show a small ( $\approx 3$ times) increase in $L_{\text {FIR }}$ compared to isolated galaxies.

Similar differences are seen in the far-IR colors. The median value of $\log \left(f_{12} / f_{25}\right)$ is -0.36 for resolved component galaxies in pairs versus -0.19 for isolated galaxies. Paper I found -0.43 and -0.2 . However, as noted earlier, this may be a result of the differing flux calibration between Paper I and this paper. Adjusting for this produces a mean $\log \left(f_{12} / f_{25}\right)$ of -0.41 for individual galaxies in the RBGS HIRES interacting galaxy sample. The maximum difference in the CDF is 0.31 and occurs at $\log \left(f_{12} / f_{25}\right)=-0.28$. Kolmogorov-Smirnov statistics reject the null hypothesis that these two samples are drawn from the same parent sample at better than the $99.99 \%$ level.

A median $\log \left(f_{60} / f_{100}\right)$ value of -0.25 is observed for galaxies belonging to pairs and groups in the RBGS HIRES sample, compared to -0.34 for noninteracting RBGS galaxies. This is the same result as seen in Paper I. The difference in CDF between the two samples is less pronounced overall than at the shorter wavelengths. Nevertheless the maximum difference in CDFs is 0.34 at $\log \left(f_{60} / f_{100}\right)=-0.29$. Again, we can reject the null hypothesis that the two samples were drawn from the same sample at better than the $99.99 \%$ level.

### 3.2. Pairing in the Far-IR

While there is clearly evidence that pairs and groups of galaxies generally have higher star formation activity compared to isolated galaxies, there is still uncertainty regarding the relative degree to which enhanced star formation is triggered in individual galaxies during various phases of the interaction and merger process. Naively, among spiral-spiral pairs that have companions with nearly equal $B$-band luminosities, one would expect that both galaxies contribute in


Fig. 8.-CDFs of $\log \left(f_{12} / f_{25}\right)$ for paired and isolated galaxies, as in Fig. 7.
similar proportions to the total far-IR emission of the pair. Previous authors, as discussed earlier, have generally concluded that for very distant pairs observed in the far-IR only one galaxy is infrared active. Other authors, working at optical and near-IR wavelengths, have reached the same conclusion using indirect measures of star formation. Using the higher resolution images presented here, it is possible to test this result over a much wider range and smaller absolute separations than previously possible. Figure 10 plots the measured flux ratios (filled circles) and upper limits (open circles) at the longest resolvable wavelength between the brightest galaxy and its companion as a function of the total far-IR luminosity, $L_{\text {FIR }}$. For galaxy groups, the ratio plotted is the flux of the brightest galaxy divided by the average flux of the companion galaxies in the group. These results provide little evidence that both companions contribute comparably to the infrared emission and that there is no increased tendency for infrared luminous galaxies to be found with other infrared luminous galaxies. That is, $66 \%$ of the interacting systems have component flux ratios greater than 3, $56 \%$ have ratios above 5 , and $36 \%$ have ratios greater than 10 . Surace et al. (1993) claimed that in approximately $\frac{2}{3}$ of interacting pairs, the ratio of the flux densities of the companions are less than 10 , which is confirmed here in the much larger sample of infrared-bright galaxy systems investigated here.

Figure 10 also shows that over the range of $L_{\text {FIR }}$ spanned by this RBGS subsample, there is no clear correlation between the companion galaxy flux ratios and $L_{\mathrm{FIR}}$. Although at flux levels $\log \left(L_{\text {FIR }} / L_{\odot}\right)>11$, suggestively there are almost no systems in which the component flux ratio is less than three, our sample does not extend to such high luminosities as would produce a clearer result because our selection criteria biases us against very advanced mergers, which are the


Fig. 9.-CDFs of $\log \left(f_{60} / f_{100}\right)$ for paired and isolated galaxies, as in Fig. 7.
majority of the luminous and ultraluminous infrared galaxies. Since both components in an interacting system are presumably undergoing a similar degree of tidal disruption, it seems that in the relatively early stages of interaction sampled here, the details of the encounter itself are less important than characteristics of the individual galaxies in determining the degree of far-IR enhancement. Major factors expected to play


Fig. 10.-Plot of $\log L_{\text {FIR }}$ vs. the pair component flux ratio. Filled circles are measured ratios, while open circles are upper limits. No correlation is seen. The distribution of lower limits is somewhat random, since there was a fixed lower flux limit for the dim component, while the bright component could span a wide range of detected brightnesses.


FIg. 11.-Integrated, spatially resolved $60 \mu \mathrm{~m}$ flux cumulative distribution functions for galaxies detected in the RBGS.
an important role are the molecular gas content available to fuel star formation and the mass of the stellar bulge, which may regulate the degree of accretion onto a supermassive black hole.

Finally, we can examine whether confusion has caused galaxies to be erroneously included in the RBGS because their combined flux was high enough to meet the flux limit criterion but which would not have been selected if they could be resolved. Figure 11 presents the cumulative distribution functions of $60 \mu \mathrm{~m}$ fluxes for the HIRES-resolved galaxies above and below the 5 Jy selection limit. An examination of Table 1 shows that there are two systems that appeared in the RBGS by virtue of having a combined, unresolved flux above the 5.24 Jy limit at $60 \mu \mathrm{~m}$ but whose individual components were all clearly below this limit (IC 563/4, NGC 7752/3). An additional five (IC 2522/3, UGC 6436a/b, NGC 3991/4/5, MCG -03-34-063a/b, and VV 414) have only one component whose flux including uncertainties may be as high as the RBGS flux limit. Thus, close pairs in the RBGS verifiably affect the selection of the sample at only the $0.3 \%$ level $(2 / 629)$ and at worst may account for $1.1 \%$. This statement applies to separated pairs resolved by HIRES in this study. Very close pairs with separations less than $\approx 30^{\prime \prime}$ (typically ongoing mergers in the local universe), which cannot be resolved by HIRES, account for increasingly larger fractions of RBGS objects as a function of increasing total far-IR luminosity (e.g., see review by Sanders \& Mirabel 1996). The manner in which the total far-IR fluxes of such objects are distributed between the individual components remains unknown for such pairs, and these will be fruitful targets to study with higher resolution using observatories such as Spitzer and SOFIA.

### 3.3. Optical Morphology and Far-IR Enhancement

Recent computational models by Mihos \& Hernquist (1994a, 1994b) have predicted that the presence of a large central bulge in a galaxy helps stabilize it against tidal perturbation. Specifically, they found that in major mergers of galaxies, the presence of a central bulge inhibits the flow of gas into the central few kiloparsecs of a galaxy, thus preventing high gas densities from being quickly reached and suppressing any period of rapid star formation until the end of the merger (Mihos \& Hernquist 1994b). Thus, late-type spirals are expected to experience starbursts during the initial stages of merger, while early-type spirals undergo strong starburst activity only during the completion of the merger process. This provides a mechanism to delay the onset of starburst activity in some systems until very advanced merger stages are reached; otherwise it is difficult to invoke a starburst model for ultraluminous infrared galaxies (which appear to be very advanced mergers) given the expected timescale for starbursts. Since very evolved mergers have such disturbed morphologies that it is difficult to determine the form of the merger progenitors, the most viable observational test is to examine young merger systems that have not evolved as far away from their original forms and look for evidence for the onset of enhanced far-IR activity in bulgeless galaxies.
Hubble types were taken from NED. A fraction of the galaxies either have not been classified at all or are simply listed with generic types such as "spiral." All other resolved spiral galaxies that were actually classified were considered to be either "early type" (S0, SB0 through Sa, and SBa) or "late type" ( $\mathrm{Sb}, \mathrm{SBb}$, and higher). Although the specific Hubble type for each galaxy was kept track of, for this analysis it was felt to be more useful to group the types into such very broad categories in order to improve the counting statistics.
Only 11 ellipticals known from optical imaging of the galaxy pairs are found in the entire sample. Of these, none are detected at both 60 and $100 \mu \mathrm{~m}$, and only three are detected by IRAS at any wavelength. These numbers agree with what would be expected based on a random pairing of elliptical and spiral galaxies given an elliptical/spiral fraction similar to that of field galaxies or are perhaps a little low. In particular, this is the number expected if we assume that every system contains one bright spiral galaxy and that the remaining faint galaxies are distributed according to the field elliptical/spiral ratio. This also agrees with the low detection fractions for elliptical galaxies found by other studies (Sulentic 1988; Haynes \& Herter 1988).

An examination of the late- and early-type spirals in the sample indicates that there are no differences between these populations. The rate of detection for both classes is around $88 \%$, indicating that they have a similar fraction of their distribution above our detection limit. This is further illustrated by Figure 12, which shows the cumulative distribution function of $\log L_{\text {FIR }}$ for the galaxies that were actually detected in the two classes. They are extremely similar. The K-S test cannot reject the null hypothesis that the two samples are drawn from the same sample with better than $65 \%$ confidence. Results are the same for the color ratios $\log \left(f_{12} / f_{25}\right)$ and $\log \left(f_{60} / f_{100}\right)$.
Similar results also hold for different combinations of Hubble subtypes, such as considering only S0 galaxies as "early." Haynes \& Herter (1988) found that the detection rate for isolated spiral galaxies is roughly independent of morphological type. Roberts \& Haynes (1994) have confirmed


FIG. 12.-Cumulative distribution function of $L_{\text {FIR }}$ of late- and early-type interacting spiral galaxies in the RBGS. There is no measurable difference between the two.
this using large optical samples of isolated galaxies and comparing them to the IRAS data. The median, $25 \%$, and $75 \%$ values of the distribution of $L_{\mathrm{FIR}}$ for the sample spiral galaxies are also very similar to those found by Roberts \& Haynes for isolated UGC galaxies. This is unsurprising, since both the results presented here and previous studies have shown that only a small increase in $L_{\mathrm{FIR}}$ occurs for widely separated (nonoverlapping) interacting pairs (Haynes \& Herter 1988; Surace et al. 1993), and therefore any enhanced far-IR activity that could distinguish late from early spirals is likely to be slight. This enhancement would be further diluted by the presence of systems that have not yet reached first perigalcticon and hence have not yet reached the point at which the two classes would separate themselves (Mihos \& Hernquist 1994b), and by the presence of systems that are unlikely to actually merge or otherwise strongly interact with each other.

## 4. CONCLUSIONS

We have presented an atlas of high-resolution IRAS observations of all 106 of the paired (and, in many cases, interacting systems) in the IRAS RBGS with a $60 \mu$ m flux density greater than 5.24 Jy . The atlas contains infrared contours overlaid on optical images and a catalog of fluxes or upper limits in all four IRAS wavebands.

We have presented the infrared luminosities and colors of the paired galaxy sample, and compared them to a sample of isolated galaxies. We find substantially the same results as Surace et al. (1993), namely, that the paired galaxies have a measurably different distribution of infrared properties than isolated galaxies.


FIG. 13.-HIRES data as in Fig. 1 for additional galaxy systems processed but not members of the final RBGS, as described in Appendix B.

Using morphological optical classifications for the galaxies, we conclude that there is no difference between late and early-type spirals in terms of their far-IR properties. In particular, no significant enhancement is seen in the far-IR luminosity or color of the late-type spirals as compared to the early-type spirals.

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## APPENDIX A

## NOTES ON INDIVIDUAL GALAXY SYSTEMS

NGC 520.- The flux in this advanced merger appears to be centered between the two galaxies, in what appears to be a dust lane.
IC 2163.-This system is somewhat puzzling in that there is significant emission to the west of the center of the western galaxy in the pair. In addition, the peak flux generally appears to occur between the galaxies. It is notable that the reconstructed IRAS beam is unfortunately elongated at nearly the same position angle as the two galaxies, and hence the irregular beam is confusing the location of the emission.

NGC 4038/4039.-The emission originates between the galaxy centers in the region where the disks overlap, which is consistent with the results of Vigroux et al. (1996).

IC 4153.-There is a sizable discrepancy between the HIRES flux and the ADDSCAN flux used by the RBGS at $100 \mu \mathrm{~m}$. Furthermore, the HIRES flux given in Table 1 agrees with the flux indicated by the FRESCO data product. FRESCO is a twodimensional coadd data product available from IPAC. This coadd is not an iterative reconstructed image and hence should not suffer from iterative artifact amplification. An examination of the complex structure seen in HIRES and FRESCO images surrounding this source, as well as the details of the ADDSCAN processing, show that this discrepancy is probably due to differences in the baseline (background) fitting.

NGC 5953/5954.-Given as CPG 468 by Domingue et al. (2003). The HIRES data agree with the higher resolution ISOPHOT data indicating the dominance of the southwestern component.

NGC 6907/6908.-The galaxy NGC 6908 is actually a small spiral galaxy superimposed on the northeast arm of NGC 6907 and is most clearly seen in near-IR images. It cannot be resolved by HIRES. This system appears in Appendix B.

NGC 7752/7753.-Also known as CPG 591. Domingue et al. (2003) find that the southwestern component is more peaked and dominates the ISOPHOT 60 and $100 \mu \mathrm{~m}$ data. The HIRES data support this finding, particularly in the mid-IR.

## APPENDIX B

## ADDITIONAL GALAXY SYSTEMS

In addition to the 106 galaxy systems detailed in $\S 2$, several other systems were also processed. During the compilation of the RBGS (Sanders et al. 2003), the IRAS data were recalibrated, and the choice of flux measures used to estimate the IRAS fluxes changed. As a result, there are a handful of systems that appear in the BGS+BGS2 (Sanders et al. 1995) but not the RBGS, and vice versa. These are detailed in the RBGS, § 3.2. As a result, some additional galaxy systems were processed with HIRES but do not belong in the RBGS sample proper and were not included in Table 1 of this paper. They are presented here for informational purposes. The images are shown in Figure 13, and the tabulated fluxes appear in Table 2.

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[^0]:    ${ }^{1}$ HIRES processing is available from the Infrared Processing and Analysis Center Infrared Science Archive, Jet Propulsion Laboratory, California Institute of Technology (http://irsa.ipac.caltech.edu/IRASdocs/hires_over.html).

